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#### *RESEARCH ARTICLE - PHYSICS*

## **Non-thermal atmospheric jet plasma effects on the morphological, electrical, and**

# **optical properties of nanocrystalline silicon**

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**Keyword: porous silicon; PECE; NTPAJ; FE-SEM; PL; XRD.**

#### **Introduction**

In the past few decades, the production of nanomaterials has become an exciting field for researchers because of their unique physical properties that differ from those of crystalline semiconductors[1]. Specifically, porous semiconductor materials, known for their distinct physical, chemical, optical, and biological properties, have received considerable attention in scientific and applied fields [2], [3] . Among the porous materials, porous silicon (PSi) is known as solid silicon with voids that provide a large surface area  $(200 \text{ to } 800 \text{ m}^2, g^{-1})$  [2], [3], nano-scale pores [4], and a combination of threedimensional semiconductor substrates with nanomaterials deposited on them, as well as a surface texture that acts as a light trap, enhancing light absorption in long wavelengths. All of these features have made it favourable for numerous applications, such as optoelectronics [5], [6], [7] biomedical [8], [9], chemical and biosensors [10], [11] and microelectronics [12]. Since the discovery of porous silicon in 1956, over 20 methods have been developed to synthesize porous silicon structures. However, only a few of these methods have gained significant attention. Out of these methods, electrochemical etching is the most widely used [13], [14], [15], [16]. It is both cost-competitive and versatile, allowing for the production of various porous structures with exceptional physical and chemical properties. By adjusting the formation parameters, such as current density, solution composition, and porous silicon etching time, a wide range of porosities and morphologies can be achieved [12], [17], [18].

Porous silicon characteristics could be improved, leading to widened applications by surface modification, which can be achieved by various methods, including chemical processes [19], annealing [20], [21], oxidation [22], as well as plasma treatment [23], [24], [25], [26]. Plasma, composed of electrons, ions, and neutral particles, is a quasi-neutral substance that can be partially or fully ionized. Thermal and non-thermal plasma are two classifications of plasma based on their component temperature, each with distinct characteristics and uses [27]. Various methods, such as plasma jet, dielectric barrier discharge, corona discharge, and arc discharge, can generate cold plasma [28], [29], [30], [31], [32], [33], [34], [35], [36], [37]. Commonly used gases include air, nitrogen, hydrogen, and argon [31], [33], [35], [37]. Non-thermal atmospheric plasma offers several advantages, including producing stable plasma at atmospheric pressure [38] and eliminating the need for a vacuum. It is also economically advantageous due to its simplicity and cost-effectiveness [37]. Additionally, nonthermal atmospheric plasma does not cause thermal damage to treated substrates [35], [38]. This type of plasma has become an essential tool in surface modification, particle deposition, sterilization of medical instruments, and various other

applications in fields like biomedical [39], insulation [40], super-wetting surfaces [41], plasma green and eco-friendly [42]. This study aims to examine how non-thermal plasma interacts with the structure of porous silicon, explicitly analyzing its impact on morphological and photoluminescence properties.

## **Experimental**

## **Photo electrochemical etching cell**

Porous silicon layers were fabricated by using the photoelectrochemical method on a n-type Si (111) substrate with a thickness of 200 µm and a resistivity of 0.001  $\Omega$ .cm. This was obtained by following a simple setup, as shown in Fig.1, and with the assistance of a halogen lamp. The substrate was cut into  $1\times1$  cm<sup>2</sup> samples, which were then thoroughly cleaned using an ultrasonic device with ethanol, acetone, and distilled water and dried with hot air. The samples were then etched for 15 minutes at a current density of 35 mA/cm<sup>2</sup>, using a solution consisting of 40% HF and 99.98% ethanol mixture with a volume ratio 1:3.



Fig. 1: Electrochemical etching cell

The porous silicon layer's structure was analyzed using x-ray diffraction with PANalytical Aeris X-ray diffractometer (Cu Kα1 radiation and wavelength of 1.54059 Å). The topography of the porous silicon surface was studied using (TT-2 AFM). Additionally, the porous silicon's morphology, thickness, and porosity were examined before and after non-thermal plasma treatment using FE-SEM (Inspect TM F50). To determine the photo luminance peak (optical energy gap), a FluoroMateFS-2 spectrometer was employed.

## **Non-thermal plasma setup**

Fig. 2 shows the designed jet consisting of two coaxial electrodes with dielectric materials, typically Teflon, positioned between them. A 20 kV AC voltage is connected to these electrodes. The inner electrode is a copper tube with a diameter of 1.5mm. In comparison, the outer electrode is a 1 cm long copper cylinder surrounding a glass tube with an inner diameter of 4mm. Argon gas flows at a rate of 10L/min within the copper tube. The plasma is generated between the electrodes and then emitted through the nozzle at the end of the glass tube onto the treatment material. Optical characterization of the discharge was done using Stark broadening and Boltzmann plot methods by analyzing the spectrum obtained using the optical emission spectrometer (HR4000CG-UV-NIR, Ocean Optics); the temperature of the electrons was estimated [43], [44], [45] and was found to be 0.547 eV.



The samples were placed at a distance of 7mm from the jet nozzle, and the treatment operation lasted for ten minutes.

### **Results and discussion**

Fig. 3 displays porous silicon's X-ray diffraction (XRD) patterns before and after argon plasma treatment. Before treatment, the samples exhibited the (111) peak at approximately 2θ = 28.39**.** After argon plasma treatment, an increase in the intensity of the (111) peak was observed, along with a slight shift towards higher angles compared to untreated porous silicon. In addition, there was an increase in the crystalline size  $(L)$  and a decrease in the average discrepancy  $(D\%)$ , as calculated using Debye-Scherrer and average discrepancy equations (Eq.1 and Eq.2) [46].

$$
L = \frac{K\lambda}{FWHM \cdot \cos\theta} \tag{1}
$$

where  $L$ : is crystalline size of porous silicon, FWHM is the full width at half maximum of the diffraction peak,  $\theta$  is the diffraction angle and  $\lambda$  is the wavelength of the Cu-K- $\alpha$ 1line.

$$
D\% = \frac{a_{si} - a_{PSi}}{a_{PSi}} \times 100\%
$$
 (2)

where D: is the lattice parameter variation between the bulk silicon constant  $a_{si}$  and the porous silicon constant  $a_{PSi}$  The full width at half maximum (FWHM) was also reduced, indicating an increase in crystalline size. These results are presented in Table 1.

Parameters	<b>Untreated Psi</b>	<b>Treated Psi</b>
$L$ (nm)	81.6 nm	288 nm
The average discrepancy D%	0.0193	0.0053
FWHM (deg.)	0.1299	0.066

Table 1: XRD result of untreated PSi and treated Psi.

For the argon plasma-treated samples, the increase in intensity and decrease in FWHM can be explained by the improvement in the crystalline structure resulting from the formation of uniform pores [47]. The strain induced by surface forces increases as the internal surface area enlarges, closely related to the porous silicon microstructure. The peak shift can be attributed to the strain, and the reduction in micro-strain can be attributed to lattice relaxation caused by changes in the surface microstructure, i.e., the size and shape of the porous silicon [20]. All of these results can be attributed to the plasma's energy rearranging the surface structure of porous silicon.



Fig. 3: XRD pattern of porous silicon layer before and after plasma treatment

The FE-SEM images of the porous silicon layer before and after argon plasma treatment (represented by figures 4 a, b, c, and d, respectively) clearly show differences in morphology. The untreated sample exhibits a semi-homogenized structure with small pentagonal and quadrangular holes forming a network. However, after plasma treatment, the walls of the pores become thicker, and the holes change their shape to nearly circular [25].

It has been observed that when porous silicon is treated with cold plasma, its porosity decreases significantly more than that of crystalline silicon that has not been treated. The plasma's energetic ions interact with the porous silicon surface to change its structural composition.

FE-SEM IMGE determined that the untreated sample's porosity was 93%, while after argon plasma treatment, the porosity decreased to 64%.



Fig. 4: FE-SEM images of a, b) As grown c, d) Plasma jet treatment of Psi

The 3D images of the porous silicon layer profile (a) and porous silicon following cold plasma treatment (b) are displayed in Figure 5. It is evident that following treatment, the semi-uniformity of pyramidal forms and surface roughness were reduced. The rearrangement and creation of pores on the surface are responsible for this roughness decrease. The amplitude reporting parameters of the samples after non-thermal plasma treatment are presented in Table 2.



Fig. 5: 3D AFM image of porous silicon a) before and b) after treatment.

Amplitude parameters	Fresh Psi	PSi /treatment
Minimum	(0.0)	(0.0)
Maximum	$108.2 \text{ nm}$	39.93 nm
(Ten Point Height) Sz	461.4 nm	$106.4 \text{ nm}$
(Roughness Average) Sa	$60.10$ nm	$12.31$ nm
(Root Mean Square) Sq	73.70 nm	$16.52 \text{ nm}$

Table 2: Roughness Report parameters of fresh PSi and PSi treated.

Fig. 6 displays the photoluminescence (PL) spectra of untreated and argon plasma-treated porous silicon at 320 nm wavelength, revealing that the PL peaks range from 560 to 680 nm for untreated porous silicon and from 520 to 645 nm for treated porous silicon. After treatment with argon plasma, specific peaks exhibited an increase in intensity and a shift towards shorter wavelengths (blue shift), as displayed at 647 and 518 nm, and some peaks vanished, as shown at 459 nm. The shift in peak position can be attributed to the luminescence mechanism of quantum confinement [48]. There are two possible reasons for the increase in intensity. One is the presence of  $SiH<sub>2</sub>$  in porous silicon and the introduction of oxygen through plasma treatment; this leads to the formation of chemical species that encourage the radiative recombination of electron-hole pairs, increasing light emission. Another possibility is that plasma will fluorinate the surface, which promotes radiative recombination and light emission [23], [26]



Fig. 6: Photoluminance spectra of porous silicon before (black line) and after (red line) treatment with atmospheric Argon plasma.

The refractive index before and after plasma treatment was calculated using the following [49]:

$$
n = \sqrt{\varepsilon_{PSi}}\n\varepsilon_{PSi} = \varepsilon_{Si} - P\%(\varepsilon_{Si} - 1)
$$
\n(3)

where  $\varepsilon_{PSi}$ : is the relative permittivity of porous silicon and equal 11.68,  $\varepsilon_{SI}$ : is the relative permittivity of bulk silicon (equal 1), and P(%): is the porosity

After plasma treatment, calculations revealed an increased refractive index of 1.32 to 2.24, so this technique may be employed to modify the optical characteristics of the porous silicon surface.

Fig. 7 shows the current -voltage characteristic for Psi/n-Si in a semi- log plot, in dark and under illumination (halogen lamp) from -8 to +8 Volts for fresh and argon plasma-treated porous silicon. The result indicated an increase in photocurrent for the non-thermal plasma treated sample, which may be attributed to enhancement in junction structure resulting from a reduction in defect at the Psi/n-Si interface, which means a decrease in the strain. The diode parameters (ideality factor n, barrier height  $\emptyset_b$ ) were calculated using the thermionic emission model equation; the relation between (I-V) is given by [50];

$$
I = I_s \left[ \exp\left(\frac{qV}{nkT}\right) - 1 \right] \tag{5}
$$

where q electronic charge, T absolute temperature in Kelvin, k Boltzmann constant, V is applied voltage, I<sub>s</sub> the reverse saturation current is calculated from the ln I vs V plot by extrapolating the linear region of the curve to zero applied voltage axis as shown in Fig. 8

$$
I_s = AA^*T^2 \exp(\frac{-q \phi_b}{nkT})
$$
  
\n
$$
\emptyset_b = \frac{kT}{q} \ln(AA^* \frac{T^2}{I_s})
$$
\n(7)

Where A is the area of the diode,  $A^*$  is the modified Richardson constant which equal to 112 A cm<sup>-2</sup> K<sup>-</sup>  $2$  for n-Si, the ideality factor is calculated from current – voltage characteristic by using the following Eq. [51]:

$$
n = \frac{q}{KT} \left( \frac{dV}{d\ln T} \right) \tag{8}
$$

The resistance  $R_d$  is calculated from the linear part of the forward bias (I-V) plot and the rectification ratio RR is calculated from the ratio between the forward to reverse current. The result obtained from the analysis of the (I-V) characteristic is summarized in Table 3.

Parameters	<b>Untreated Psi</b>	<b>Treated Psi</b>
$\mathbf{1}_{\mathbf{S}}$	$5.25 \times 10^{-7}$	$6.14\times10^{-6}$
n	2.4	2.6
$\emptyset_b$	0.758	0.683
$R_{d}$	$167 \text{ k}\Omega$	$100 \text{ k}\Omega$
<b>RR</b>		33

Table 3: (I-V) characteristic of fresh PSi and PSi treated.

The results indicated that after argon plasma treatment, there is an increase in reverse saturation current accompanied by a decrease in the barrier height and a decrease in resistance, which may be attributed to the reduction in layer thickness and porosity. A non-homogeneously porous silicon layer may lead to a high ideality factor of more than one, and the increase in ideality factor may be caused by the surface oxidation of porous silicon after non-thermal plasma treatment.

The increased photo-current can be attributed to several factors. Firstly, cold plasma treatment can enhance the surface passivation of porous silicon, and secondly, modify the surface properties of porous silicon, leading to changes in its electrical behaviour and reducing the surface recombination of photo-generated carriers. This reduction in carrier recombination improves the overall efficiency of the material, leading to higher photo-current generation.



Fig. 7: Photocurrent of PSi/n-Si diode before (red line) and after (blue line) plasma treatment.



## **Conclusions**

This study investigated the impact of jet plasma on the morphological, electrical, and optical properties of nanocrystalline porous silicon. The photoelectrochemical etching method was used to fabricate a porous silicon layer.

The n-type porous silicon surface was effectively created, and its microstructure, morphological, and electrical properties were compared to those of the untreated surface.The surface of porous silicon can be passivity by non-thermal plasma by forming a barrier. The electrical characteristics of porous silicon devices, like solar cells or sensors, can be improved, and this passivation layer can decrease surface recombination. It contributes to increased device stability and efficiency.Low-temperature plasma treatment is beneficial because it maintains the material's structural integrity and porosity, which are essential for its special qualities and uses. Treating the PSi surface with cold plasma could further improve the morphological, optical, electrical, and optoelectronic capabilities of the porous silicon layer, it contributes to increased device stability and efficiency.Modifying and treating nanostructures of porous silicon using non-thermal jet plasma technologies are very promising.

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