

**Extending Application of Adomian Decomposition Method for  
Solving a Class of Volterra Integro-Differential Equations  
within Local Fractional Integral Operators**

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**ABSTRACT**

In this paper, we consider the local fractional Adomian decomposition method for solving the second kind Volterra integro-differential equations within local fractional integral operators. This application maintains the efficiency and accuracy of the Adomian analytic method for solving local fractional integral equations. Illustrative examples are given to show the accuracy and reliability of the results.

**Keywords:** Adomian decomposition method; Local fractional integral operators; Local fractional Volterra integro-differential equations.

توسيع تطبيق طريقة تحليل ادوميان لحل صنف من معادلات فولترا التكاملية التفاضلية

ضمن المؤثرات التكاملية الكسرية المحلية

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#### الخلاصة

في هذا البحث ندرس طريقة تحليل ادوميان الكسرية المحلية لحل معادلات فولترا التكاملية التفاضلية من النوع الثاني ذات المؤثر التكاملي الكسري المحلي. هذا التطبيق يؤكد كفاءة ودقة طريقة ادوميان التحليلية لحل المعادلات التكاملية الكسرية المحلية. قدمت امثلة توضيحية تبين دقة النتائج.

كلمات البحث: طريقة ادوميان التحليلية, المؤثرات التكاملية الكسرية المحلية, معادلات فولترا التكاملية التفاضلية.

## 1. Introduction:

Local fractional calculus is a new branch of mathematics that deals with derivatives and integrals of the functions defined on fractal sets. It is used to explain behavior of continuous but nowhere differentiable function. Local fractional calculus (is also called Fractal calculus) was first introduced by Kolwankar and Gangal [1].

Recently, the local fractional calculus theory has attracted a lot of interest for scientists and engineers because it is applied to model problems in fractal mathematics and engineering. It plays a key role in many applications in several fields, such as physics [2,3], applied mathematics [4,5], signal processing [6,7], fluid mechanics [8], quantum mechanics [9], fractal forest gap [10], vehicular traffic flow [11], and silk cocoon hierarchy [12]. There are many analytical and numerical methods used to solve differential equations on Cantor sets such as, Adomian decomposition method [13,14], variational iteration method [15-17], series expansion method [18,19], homotopy perturbation method [20,21], Laplace Variational Iteration Method [22,23], differential transform method [24,25], function decomposition method [26], and similarity solution [27] involving the local fractional operators.

In this work, we consider analysis solution to the local fractional Volterra integro-differential equation of the second kind. This paper is organized as follows: In section 2, the concept of local fractional calculus and integrals are given. In section 3, Adomian decomposition method is proposed based on local fractional integrals. Illustrative examples are shown in section 4. Conclusions are given in section 5.

## 2. Preliminary Definitions

**Definition 1** [1,13]: The function  $f(x)$  is local fractional continuous at  $x = x_0$ , if it is valid for

$$|f(x) - f(x_0)| < \varepsilon^\alpha, 0 < \alpha \leq 1 \quad (2.1)$$

with  $|x - x_0| < \delta$ , for  $\varepsilon > 0$  and  $\varepsilon \in R$ . For  $x \in (a, b)$ ; it is so called local fractional continuous on the interval  $(a, b)$ , denoted by  $f(x) \in C_\alpha(a, b)$ .

**Definition 2**[13]: The local fractional derivative of  $f(x)$  at  $x = x_0$  is defined as

$$D_x^\alpha f(x_0) = \frac{d^\alpha}{dx^\alpha} f(x) \Big|_{x=x_0} = f^{(\alpha)}(x) = \lim_{x \rightarrow x_0} \frac{\Delta^\alpha (f(x) - f(x_0))}{(x - x_0)^\alpha} \quad (2.2)$$

where  $\Delta^\alpha (f(x) - f(x_0)) \cong \Gamma(\alpha + 1) \Delta(f(x) - f(x_0))$ .

**Definition 3**[13]: A partition of the interval  $[a,b]$  is denoted as  $(t_j, t_{j+1})$ ,  $j=0, \dots, N-1$ ,  $t_0 = a$  and  $t_N = b$  with  $\Delta t_j = t_{j+1} - t_j$  and  $\Delta t = \max\{\Delta t_0, \Delta t_1, \dots\}$ . The local fractional integral of  $f(x)$  in the interval  $[a,b]$  is given by

$${}_a I_b^{(\alpha)} f(x) = \frac{1}{\Gamma(1+\alpha)} \int_a^b f(t) (dt)^\alpha = \frac{1}{\Gamma(1+\alpha)} \lim_{\Delta t \rightarrow 0} \sum_{j=0}^{N-1} f(t_j) (\Delta t_j)^\alpha. \quad (2.3)$$

**Definition 4** [13]: In fractal space, the MittagLeffler function, sine function, cosine function are, respectively defined by

$$E_\alpha(x^\alpha) = \sum_{k=0}^{\infty} \frac{x^{k\alpha}}{\Gamma(1+k\alpha)}, \quad 0 < \alpha \leq 1 \quad (2.4)$$

$$\sin_\alpha(x^\alpha) = \sum_{k=0}^{\infty} (-1)^k \frac{x^{(2k+1)\alpha}}{\Gamma[1+(2k+1)\alpha]}, \quad 0 < \alpha \leq 1 \quad (2.5)$$

$$\cos_\alpha(x^\alpha) = \sum_{k=0}^{\infty} (-1)^k \frac{x^{2k\alpha}}{\Gamma[1+2k\alpha]}, \quad 0 < \alpha \leq 1 \quad (2.6)$$

### 3. Analysis of the Method

The standard  $k\alpha$  order local fractional Volterra integro-differential equation of the second kind is given by

$$u^{(k\alpha)}(x) = f(x) + \frac{\lambda^\alpha}{\Gamma(1+\alpha)} \int_0^x K(x,t) u(t) (dt)^\alpha, \quad (3.1)$$

subject to initial conditions

$$u(0) = a_0, u^{(\alpha)}(0) = a_1, \dots, u^{((k-1)\alpha)}(0) = a_{k-1},$$

where  $u^{(k\alpha)}(x) = \frac{d^{k\alpha} u(x)}{dx^{k\alpha}}$  is linear local fractional derivative operator.

By integrating both sides of (3.1) leads to

$$L^{(-k\alpha)}(u^{(k\alpha)}(x)) = L^{(-k\alpha)}(f(x)) + L^{(-k\alpha)}\left(\frac{\lambda^\alpha}{\Gamma(1+\alpha)} \int_0^x K(x,t) u(t) (dt)^\alpha\right) \quad (3.2)$$

Thus, we obtain

$$u(x) = a_0 + a_1 \frac{x^\alpha}{\Gamma(1+\alpha)} + \dots + a_{k-1} \frac{x^{(k-1)\alpha}}{\Gamma[1+(k-1)\alpha]} + L^{(-k\alpha)}(f(x)) + L^{(-k\alpha)}\left(\frac{\lambda^\alpha}{\Gamma(1+\alpha)} \int_0^x K(x,t) u(t) (dt)^\alpha\right) \quad (3.3)$$

where the initial conditions  $u(0), u^{(\alpha)}(0), u^{(2\alpha)}(0), \dots, u^{((k-1)\alpha)}(0)$  are used, and

$$L^{(-k\alpha)}(\cdot) = {}_0 I_x^{(k\alpha)} = \overbrace{{}_0 I_x^{(\alpha)} {}_0 I_x^{(\alpha)} \dots {}_0 I_x^{(\alpha)}}^{k \text{ time}}(\cdot). \tag{3.4}$$

Then, we use the decomposition series

$$u(x) = \sum_{n=0}^{\infty} u_n(x) \tag{3.5}$$

in both sides (3.3) to obtain

$$\sum_{n=0}^{\infty} u_n(x) = a_0 + a_1 \frac{x^\alpha}{\Gamma(1+\alpha)} + \dots + a_{k-1} \frac{x^{(k-1)\alpha}}{\Gamma(1+(k-1)\alpha)} + L^{(-k\alpha)}(f(x)) + L^{(-k\alpha)}\left(\frac{\lambda^\alpha}{\Gamma(1+\alpha)} \int_0^x K(x,t) \left(\sum_{n=0}^{\infty} u_n(t)\right) (dt)^\alpha\right) \tag{3.6}$$

or equivalently

$$u_0(x) + u_1(x) + u_2(x) + \dots = a_0 + a_1 \frac{x^\alpha}{\Gamma(1+\alpha)} + \dots + a_{k-1} \frac{x^{(k-1)\alpha}}{\Gamma(1+(k-1)\alpha)} + L^{(-k\alpha)}(f(x)) + L^{(-k\alpha)}\left(\frac{\lambda^\alpha}{\Gamma(1+\alpha)} \int_0^x K(x,t) u_0(t) (dt)^\alpha\right) + L^{(-k\alpha)}\left(\frac{\lambda^\alpha}{\Gamma(1+\alpha)} \int_0^x K(x,t) u_1(t) (dt)^\alpha\right) + \dots \tag{3.7}$$

To determine the components  $u_0(x), u_1(x), u_2(x), \dots$  of the solution  $u(x)$  we set the recurrence relations

$$u_0(x) = a_0 + a_1 \frac{x^\alpha}{\Gamma(1+\alpha)} + \dots + a_{k-1} \frac{x^{(k-1)\alpha}}{\Gamma(1+(k-1)\alpha)} + L^{(-k\alpha)}(f(x)) \tag{3.8}$$

$$u_{n+1}(x) = L^{(-k\alpha)}\left(\frac{\lambda^\alpha}{\Gamma(1+\alpha)} \int_0^x K(x,t) u_n(t) (dt)^\alpha\right), \quad n \geq 0. \tag{3.9}$$

It means that

$$u_0(x) = a_0 + a_1 \frac{x^\alpha}{\Gamma(1+\alpha)} + \dots + a_{k-1} \frac{x^{(k-1)\alpha}}{\Gamma(1+\alpha)} + L^{(-k\alpha)}(f(x)) \tag{3.10}$$

$$u_1(x) = L^{(-k\alpha)}\left(\frac{\lambda^\alpha}{\Gamma(1+\alpha)} \int_0^x K(x,t) u_0(t) (dt)^\alpha\right). \tag{3.11}$$

$$u_2(x) = L^{(-k\alpha)}\left(\frac{\lambda^\alpha}{\Gamma(1+\alpha)} \int_0^x K(x,t) u_1(t) (dt)^\alpha\right). \tag{3.12}$$

$$u_3(x) = L^{(-k\alpha)}\left(\frac{\lambda^\alpha}{\Gamma(1+\alpha)} \int_0^x K(x,t) u_2(t) (dt)^\alpha\right). \tag{3.13}$$

#### 4. An Illustrative Examples

In this section three examples for the local fractional Volterra integro-differential equation from the second kind is presented in order to demonstrate the simplicity and the efficiency of the above method.

**Example 1.** We consider the local fractional Volterra integro-differential equation

$$u^{(\alpha)}(x) = 1 - \frac{1}{\Gamma(1+\alpha)} \int_0^x u(t)(dt)^\alpha, u(0) = 0 \quad (4.1)$$

Let the solution in the series form

$$u(x) = \sum_{n=0}^{\infty} u_n(x) \quad (4.2)$$

Applying the integral operator  $L^{(-\alpha)}$  defined by

$$L^{(-\alpha)}(\cdot) = I_x^{(\alpha)}(\cdot) = \frac{1}{\Gamma(1+\alpha)} \int_0^x (\cdot)(dt)^\alpha \quad (4.3)$$

to both sides of (4.1), and using the given initial condition we obtain

$$u(x) = \frac{x^\alpha}{\Gamma(1+\alpha)} - L^{(-\alpha)} \left( \frac{1}{\Gamma(1+\alpha)} \int_0^x u(t)(dt)^\alpha \right) \quad (4.4)$$

Then substituting (4.2) in (4.4), we have that

$$\sum_{n=0}^{\infty} u_n(x) = \frac{x^\alpha}{\Gamma(1+\alpha)} - L^{(-\alpha)} \left( \frac{1}{\Gamma(1+\alpha)} \int_0^x \left( \sum_{n=0}^{\infty} u_n(t) \right) (dt)^\alpha \right) \quad (4.5)$$

From (3.8) and (3.9), we get

$$u_0(x) = \frac{x^\alpha}{\Gamma(1+\alpha)} \quad (4.6)$$

$$u_1(x) = L^{(-\alpha)} \left( \frac{1}{\Gamma(1+\alpha)} \int_0^x u_0(t)(dt)^\alpha \right) = -\frac{x^{3\alpha}}{\Gamma(1+3\alpha)} \quad (4.7)$$

$$u_2(x) = L^{(-\alpha)} \left( \frac{1}{\Gamma(1+\alpha)} \int_0^x u_1(t)(dt)^\alpha \right) = \frac{x^{5\alpha}}{\Gamma(1+5\alpha)} \quad (4.8)$$

$$u_3(x) = L^{(-\alpha)} \left( \frac{1}{\Gamma(1+\alpha)} \int_0^x u_2(t)(dt)^\alpha \right) = -\frac{x^{7\alpha}}{\Gamma(1+7\alpha)} \quad (4.9)$$

and so on.

This gives the solution in a series form

$$u(x) = \frac{x^\alpha}{\Gamma(1+\alpha)} - \frac{x^{3\alpha}}{\Gamma(1+3\alpha)} + \frac{x^{5\alpha}}{\Gamma(1+5\alpha)} - \frac{x^{7\alpha}}{\Gamma(1+7\alpha)} \dots$$

$$= \sum_{k=0}^{\infty} (-1)^k \frac{x^{(2k+1)\alpha}}{\Gamma[1+(2k+1)\alpha]} \quad (4.10)$$

and hence the exact solution is given by

$$u(x) = \sin_\alpha(x^\alpha) \quad (4.11)$$

**Example 2.** We consider the local fractional Volterra integro-differential equation

$$u^{(2\alpha)}(x) = 1 + \frac{x^\alpha}{\Gamma(1+\alpha)} + \frac{1}{\Gamma(1+\alpha)} \int_0^x \frac{(x-t)^\alpha}{\Gamma(1+\alpha)} u(t) (dt)^\alpha, u(0) = 1, u^{(\alpha)}(0) = 1 \quad (4.12)$$

Let the solution in the series form

$$u(x) = \sum_{n=0}^{\infty} u_n(x) \quad (4.13)$$

Applying the integral operator  $L^{(-2\alpha)}$  to both sides of (4.12), and using the given initial condition we obtain

$$u(x) = 1 + \frac{x^\alpha}{\Gamma(1+\alpha)} + \frac{x^{2\alpha}}{\Gamma(1+2\alpha)} + \frac{x^{3\alpha}}{\Gamma(1+3\alpha)} + L^{(-2\alpha)} \left( \frac{1}{\Gamma(1+\alpha)} \int_0^x \frac{(x-t)^\alpha}{\Gamma(1+\alpha)} u(t) (dt)^\alpha \right) \quad (4.14)$$

Then substituting (4.13) in (4.14), we have that

$$\sum_{n=0}^{\infty} u_n(x) = 1 + \frac{x^\alpha}{\Gamma(1+\alpha)} + \frac{x^{2\alpha}}{\Gamma(1+2\alpha)} + \frac{x^{3\alpha}}{\Gamma(1+3\alpha)} + L^{(-2\alpha)} \left( \frac{1}{\Gamma(1+\alpha)} \int_0^x \frac{(x-t)^\alpha}{\Gamma(1+\alpha)} \left( \sum_{n=1}^{\infty} u_n(t) \right) (dt)^\alpha \right)$$

$$u_0(x) = a_0 + a_1 \frac{x^\alpha}{\Gamma(1+\alpha)} + L^{(-k\alpha)}(f(x)) \quad (4.15)$$

$$u_{n+1}(x) = L^{(-k\alpha)} \left( \frac{1}{\Gamma(1+\alpha)} \int_0^x \frac{(x-t)^\alpha}{\Gamma(1+\alpha)} u_n(t) (dt)^\alpha \right), n \geq 0. \quad (4.16)$$

Therefore,

$$u_0(x) = 1 + \frac{x^\alpha}{\Gamma(1+\alpha)} + \frac{x^{2\alpha}}{\Gamma(1+2\alpha)} + \frac{x^{3\alpha}}{\Gamma(1+3\alpha)} \quad (4.17)$$

$$u_1(x) = L^{(-2\alpha)} \left( \frac{1}{\Gamma(1+\alpha)} \int_0^x \frac{(x-t)^\alpha}{\Gamma(1+\alpha)} u_0(t) (dt)^\alpha \right)$$

$$= \frac{x^{4\alpha}}{\Gamma(1+4\alpha)} + \frac{x^{5\alpha}}{\Gamma(1+5\alpha)} + \frac{x^{6\alpha}}{\Gamma(1+6\alpha)} + \frac{x^{7\alpha}}{\Gamma(1+7\alpha)} \quad (4.18)$$

and so on. This gives the solution in a series form

$$\begin{aligned}
 u(x) &= 1 + \frac{x^\alpha}{\Gamma(1+\alpha)} + \frac{x^{2\alpha}}{\Gamma(1+2\alpha)} + \frac{x^{3\alpha}}{\Gamma(1+3\alpha)} + \frac{x^{4\alpha}}{\Gamma(1+4\alpha)} + \frac{x^{5\alpha}}{\Gamma(1+5\alpha)} + \dots \\
 &= \sum_{k=0}^{\infty} \frac{x^{k\alpha}}{\Gamma(1+k\alpha)}
 \end{aligned}
 \tag{4.19}$$

Hence, the exact solution

$$u(x) = E_\alpha(x^\alpha). \tag{4.20}$$

**Example 3.** We consider the local fractional Volterra integro-differential equation

$$\begin{aligned}
 u^{(4\alpha)}(x) &= -1 + \frac{x^\alpha}{\Gamma(1+\alpha)} - \frac{1}{\Gamma(1+\alpha)} \int_0^x \frac{(x-t)^\alpha}{\Gamma(1+\alpha)} u(t) (dt)^\alpha \\
 u(0) &= -1, u^{(\alpha)}(0) = 1, u^{(2\alpha)}(0) = 1, u^{(3\alpha)}(0) = -1
 \end{aligned}
 \tag{4.21}$$

Applying Equations (3.8) and (3.9), we arrive at the following iteration formula:

$$u_0(x) = -1 + \frac{x^\alpha}{\Gamma(1+\alpha)} + \frac{x^{2\alpha}}{\Gamma(1+2\alpha)} - \frac{x^{3\alpha}}{\Gamma(1+3\alpha)} + L^{(-4\alpha)}(f(x)) \tag{4.22}$$

$$u_{n+1}(x) = -L^{(-k\alpha)} \left( \frac{1}{\Gamma(1+\alpha)} \int_0^x \frac{(x-t)^\alpha}{\Gamma(1+\alpha)} u_n(t) (dt)^\alpha \right), n \geq 0. \tag{4.23}$$

$$u_0(x) = -1 + \frac{x^\alpha}{\Gamma(1+\alpha)} + \frac{x^{2\alpha}}{\Gamma(1+2\alpha)} - \frac{x^{3\alpha}}{\Gamma(1+3\alpha)} - \frac{x^{4\alpha}}{\Gamma(1+4\alpha)} + \frac{x^{5\alpha}}{\Gamma(1+5\alpha)} \tag{4.24}$$

$$\begin{aligned}
 u_1(x) &= -L^{(-4\alpha)} \left( \frac{1}{\Gamma(1+\alpha)} \int_0^x \frac{(x-t)^\alpha}{\Gamma(1+\alpha)} u_0(t) (dt)^\alpha \right) \\
 &= \frac{x^{6\alpha}}{\Gamma(1+6\alpha)} + \frac{x^{7\alpha}}{\Gamma(1+7\alpha)} + \frac{x^{8\alpha}}{\Gamma(1+8\alpha)} + \dots
 \end{aligned}
 \tag{4.25}$$

and so on.

This gives the solution in a series form

$$\begin{aligned}
 u(x) &= \left( \frac{x^\alpha}{\Gamma(1+\alpha)} - \frac{x^{3\alpha}}{\Gamma(1+3\alpha)} + \frac{x^{5\alpha}}{\Gamma(1+5\alpha)} \dots \right) - \left( 1 - \frac{x^{2\alpha}}{\Gamma(1+2\alpha)} + \frac{x^{4\alpha}}{\Gamma(1+4\alpha)} + \dots \right) \\
 &= \sin_\alpha(x^\alpha) - \cos_\alpha(x^\alpha).
 \end{aligned}
 \tag{4.26}$$

### 5. Conclusions

In this work, we present the analytical approximation to a solution local fractional Volterra integro-differential equation. We have achieved this goal by applying Adomian decomposition method. Using the ADM, it is possible to find the exact solution or a good approximate solution of the



equation. It can be concluded that ADM is a very powerful and efficient technique for finding exact solutions for wide classes of problems.

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