Compute Fermi hole and the<u>Expectation</u> values for Li-atom Using DZ, SZ, Muilt-ZateData

DZ,SZ, حساب فجوة فيرمي والقيم المتوقعة لذرة الليثيوم باستخدام بيانات Muilt-Zate

Entessar Farhan Selman Kufa University, College of scince, Department of physics. fentesser@yahoo.com

Abstract

In this paperthe partitioning method has been used for the atomic system consist of threeelectrons (Li-atom), Hartree-Fock energy for Li-atom have been calculated using the Hartree-Fock wave functions of Clementi and Rotti by a single orbital "Single Zate" and "Double Zate", also it calculated using Hatree-Fock wave functions for Mulit-Zate by Sarsa *et al*, these wave functions were found to give rather inaccurate values to DZ for Clementi and Rotti, but indicating some " exact" total Hartree-Fock energy as a functions of the Single-Zate and 7-term. The expectation values of $\langle r_1^k \rangle$, $\langle r_{12}^k \rangle$, Δr_1 , Δr_{12} , $D(r_1)$, $f(r_{12})$, $\langle V_{ee} \rangle$, $\langle V \rangle$, $\langle E \rangle$ and Fermi hole considered by DZ, SZ, mulit-Zate. The differences between intra-and inter-shell correlation estimation using the different wave functions is made.

الخلاصة

في هذه البحث استخدمت طريقة التجزئة للنظام الذري الذي يتألف من ثلاثة الكترونات (ذرة الليثيوم). طاقة هارتري-فوك لذرة الليثيوم حسبت باستخدام الدوال الموجية لمهارتري-فوك لروتي وكلمنتي للاوربيتال المفرد "Single Zate , والاوربيتال التنائي Double-Zate كذلك حسبت باستخدام الدوال الموجية لمهارتري – فوك Mulit-Zate للسارسا واخرون وجد ان هذه الدوال تعطي في الواقع قيم غير مضبوطة للطاقة , DZ, ولكن تضمنت الطاقة الكلية لمهارتري-فوك واخرون وجد ان هذه الدوال تعطي في الواقع قيم غير مضبوطة للطاقة , p_{2} , $\langle V_{2}$,

1.Introduction

The most important appliance of Har-tree-Fock (HF) method was, in atomic and molecular physics, the study of Coulomb systems (atomic, ions and molecular) with a entirely Coulombic Hamiltonian of electrons interacting with static nuclei. In physics, the use of Hartree-Fock method to calculate the energy of ground state of atoms, ions and molecular, as well compute the atomic prop-erties we refer, for example to papers[1,2,3,4]. Roothaan's open shell HF theory[5] is universally used to review the diverse of defi-nite states of atoms, molecular and nuclei .

In this paper we used the partitioning technique has been used to analyze Li-atom for six-pairs electronic wave functions, two of these are for the intra-shells (K-and L-shells) and the rest for inter-shells by using Hartree-Fock wave functions published by Clementi and Roetti for single-and double Zate[6], and we used this method for study the Roothaan Hartree-Fock wave functions published by Sarsa, Gálvez and Buendía in paper them[7] for 7-term.

In this report we focus our work on the determination of the trial parameters for improved versions of ground state $({}^{1}S, {}^{3}S)$ wave functions, for Li-atom.Mean energies, expectation value for distances for r_{i} , r_{ij} (electron-nucleus, electron-electron) respectively and Fermi hole are tabulated for tai-lored wave functions here posed. Atomic units are used unless otherwise stated.

2. Theory:

The orbital effective potential method which using in this paper have been obtained by [7] is a variational approximation to the many-electron problem which uses single-particle wave functions as variational ansatz as in Hartree-Fock mehod. The single particle orbitals satisfy a single-particle Schrödinger equation with a certain local potential that is taken to be central [8],

$$\left(-\frac{1}{2}\vec{\nabla}^2 + V(r)\right)\phi_{\lambda\sigma}(\vec{r}) = \varepsilon_{\lambda}\phi_{\lambda\sigma}(\vec{r}), \quad (1)$$

Where $\lambda = (n, l, m_l)$ stands for the spatial quantum numbers of atomic orbital and $\sigma = \pm 1/2$ is the third component of the electron spin, ∇_i^2 the Laplace operator, ε_{λ} is energy and given by

$$\varepsilon_{\lambda} = I_{\lambda} + \sum_{\mu} (J_{\lambda\mu} - K_{\lambda\mu}),$$
 (2)

I, *J* and *K* are usuall single-particle, direct and exchange integrals calculated starting from the single-particle orbitals. The wave function $\psi_{HF}(r_1, r_2, ..., r_N)$ is a Slater determinant constructed from the orbital ϕ_i [9],

$$\psi_{HF}(r_1, r_2, \dots, r_N) = A(\emptyset_1, \dots, \emptyset_N)$$
(3)

Where A is the antisymmetrizer and N is the total number of electrons. The orbitals are expanded in terms of a set of basis functions[10],

$$\phi_i(r) = \sum_i C_{ij} x_{ij}(r) \tag{4}$$

Where the basis functions $x_{ij}(r)$ are Slatertype orbitals with integer quantum numbers. The expanded coefficients C_{ij} and orbital exponents(denoted ξ_i) for each Slater-type orbital are provided in tables by Clementi and Roetti. The basis functions $x_{ij}(r)$ are slater-type orbitals with integer quantum number, namely[11]

$$x_{ij}(r) = R(r) (5)$$

where $Y(\theta, \phi)$ are normalized spherical harmonics R(r) is the radial part[12],

$$R(r) = [(2n)!]^{-1/2} (2\xi)^{n+1/2} r^{n-1} e^{-\xi r}$$
(6)

The exponent ξ , and principle quantum numbers,n,of the wave functions of Clementi and Roetti based on Slater-type function(STF_s) and showed that the optimum choice of n improves notably the double-zeta(DZ) approximation of atoms. A double-Zeta function is an approximate Roothaan-Hartree-Fock function in which a given electron orbitals is described by two Slater functions. While single-zeta function is described by one Slater functions. And milt-zeta function is described by mulit Slater functions[6]. We calculated the one-particle expectation value by equation[13]

$$< r_1^n > = \int_0^\infty D(r_1) r_1^n dr_1$$
 (7)

Where $D(r_l)$ is the one particle density distribution function. The two-particle expectation values calculated from equation[14],

$$\langle r_{ij}^n \rangle = \int_0^\infty f(r_{ij}) r_{ij}^n dr_{ij}$$
(8)

 $f(r_{ij})$, interparticle radial density distribution function, we determined, fermi holeFermi hole $\Delta f_F(r_{12})$ is defined as a difference between the electron-electron distribution function $f(r_{12})$ for triplet state and these for singlet state[15],

$$\Delta f_F(r_{12}) = f_{\alpha\alpha}(r_{12}) - f_{\alpha\beta}(r_{12}) \qquad (9)$$

3. Results and discussion

Tables 1, 2 summarizer the RHF one expectation values, two-expectation values respectively, for Li atom obtained from SZ and DZ wave functions published by Clementi and Rotti and muilt- Zeta wave function published by Sarsa *et al*(2004).

From Table 1 and Table 2, we showed that the results of the one-particle expectation values, twoparticle expectation values and standard deviations Δr_1 , Δr_{12} by Single-exponent ξ , and single principle quantum number agree with mult-exponent ξ and with Ref[7] with smallest different, comparison with Double-exponent ξ .

Shell	$\langle r_1^{-2} \rangle$	$\langle r_1^{-1} \rangle$	$\langle r_1^1 \rangle$	$\langle r_1^2 \rangle$	Δr_1
$k_{\alpha}k_{\betaSZ}$	14.40451	2.68048	0.55988	0.42219	0.32974
$k_{\alpha}k_{\betaDZ}$	15.45881	2.69024	0.54265	0.41380	2.11295
muilt-term $k_{\alpha}k_{\beta}$	14.88317	2.68386	0.57361	0.44830	0.34534
<i>Ref</i> [16]	14.8883	2.6850	0.5731	0.4468	
$K_{\alpha}L_{\alpha}(1s)_{SZ}$	4.52714	1.53041	2.28267	9.67073	2.11191
$K_{\alpha}L_{\alpha}(1s)_{DZ}$	6.24240	1.04327	0.50560	0.85385	0.77345
$K_{\alpha}L_{\alpha}(1s)_{muilt-term}$	7.67179	1.51575	2.22188	9.05982	2.03053
<i>Ref</i> [17]			2.2231	9.0874	
$K_{\beta}L_{\alpha}(3s)_{SZ}$	7.59728	1.54000	2.27913	9.65896	2.11295
$K_{\beta}L_{\alpha} (3s)_{DZ}$	7.59728	1.53996	2.27912	9.65896	2.11295
$K_{\beta}L_{\alpha}(3s)_{muilt-term}$	7.67179	1.51575	2.221188	9.05982	2.11295
<i>Ref</i> [18]	7.662	1.515	2.223	9.088	

 Table 1: Comparison of the one-particle expectation values for Li atom obtained from wave functions Single- Double-Zeta(RHF) by Clementi and Rotti with wave functions (RHF) obtained from 7-term by Sarsa *et al.*

*Note energies are expressed in Atomic units.

Shell	$\langle r_{12}^{-2} \rangle$	$\langle r_{12}^{-1} angle$	$\langle r_{12}^1 angle$	$\langle r_{12}^2 \rangle$	Δr_{12}
$k_{\alpha}k_{\beta} SZ$	4.79282	1.67353	0.81718	0.84439	0.42024
$k_{\alpha}k_{\beta DZ}$	0.24171	0.33328	4.03955	19.31792	1.73203
muilt-term $k_{\alpha}k_{\beta}$	4.71247	1.64896	0.84032	0.89659	0.43640
<i>Ref</i> [18]	4.7172	1.6498	0.8393	0.8936	0.4346
$K_{\alpha}L_{\alpha}(1s)_{SZ}$	0.12288	0.30951	4.05158	19.34146	1.71060
$K_{\alpha}L_{\alpha}(1s)_{DZ}$	0.33266	0.35686	0.82803	17.0771	1.01098
$K_{\alpha}L_{\alpha}(1s)_{muilt-term}$	0.12013	0.30841	3.92457	18.11964	1.64845
<i>Ref</i> [19,17]	0.1901	0.3227	3.9221	18.1848	
$K_{\beta}L_{\alpha}(3s)_{SZ}$	0.24171	0.33328	4.03955	19.31792	1.73202
$K_{\beta}L_{\alpha}(3s)_{DZ}$	0.24171	0.33328	4.03955	19.31792	1.73202
$K_{\beta}L_{\alpha}$ (3s) muilt-term	0.19020	0.32341	3.91892	18.11964	1.66184
<i>Ref</i> [19,17]	0.1199	0.3084	3.9274	18.1848	

 Table 2: Comparison of the two-particle expectation values for Li atom obtained from wave functions Single- Double-Zeta(RHF) by Clementi and Rotti with wave functions (RHF) obtained from 7-term by Sarsa *et al.*

Table 3: Expectation values of the potential energy $\langle V \rangle$ and total energy $\langle E \rangle$ of the HF wave function for Li-atom for data(1974) [6], and data forLi-atom(2004) [7].

Shell	$-V_{en} = -Z^* < r^{-1} >$	V _{ee}	$-\langle V \rangle$	$\langle -E \rangle$
$k_{\alpha}k_{\beta} SZ$	8.04144	1.67353	6.36791	3.18396
$K_{\alpha}L_{\alpha}(1s)_{SZ}$	4.59123	0.30951	4.28172	2.14086
$K_{\beta}L_{\alpha}(3s)_{SZ}$	3.03069	0.33328	4.25700	2.12850
Total energy(SZ)		••••	14.90663	7.45332
$k_{\alpha}k_{\beta} DZ$	8.07072	0.33328	7.73744	3.86872
$K_{\alpha}L_{\alpha}(1s)_{DZ}$	3.12981	0.35686	2.77295	1.38647
$K_{\beta}L_{\alpha}(3s)_{DZ}$	4.61988	0.33328	4.28660	2.14330
Total energy(DZ)			14.79699	7.39850
$k_{\alpha}k_{\beta}$ muilt-term	8.05158	2.68386	6.40262	3.20131
$K_{\alpha}L_{\alpha}(1s)_{muilt-term}$	4.54725	0.30841	4.23884	2.11942
$K_{\beta}L_{\alpha}$ (3s) muilt-term	4.54725	0.30841	4.23884	2.11942
Total energy _{muilt-term}			14.86530	7.43265
Experimental[20]				7.47807

In table 3, evaluated the repulsion energy expectation values $\langle V_{ee} \rangle$, the attractive energy expectation values $\langle V_{en} \rangle$ and total energy expectation values $\langle E \rangle$ for Li-atom by using expectation values for $\langle r^{-1} \rangle$ and $\langle r_{12}^{-1} \rangle$ fromtables 1,2. We can concluded the total energy expectation value obtained by two-Slater (DZ) determinant is good agreement with multi-Slater(mult-term) determinant with experimental value (-7.279913) but the total energy expectation value obtained by one-Slater(SZ) inaccurate. This paper has been designed to investigate the physical properties reflected from the figures 1,2. The result which produced the main atomic properties are one- particle density $D(r_1)$ and inter particle distribution function $f(r_{12})$ One may find the following notes For double-Zeta orbital including those resulting errors in singlet and triple states where KL(${}^{1}S$) \neq KL(${}^{3}S$), but given some accurate for both Single-, mult-zate wave functions.





Figure(1): The one particle radial density function D(r1) for Li-atom



7-term –Zate

Figure (2) inter-particle distribution function f(r12) for Li-atom.









double -Zate

Figure 3: Fermi hole is plotted as a difference between $f(r_{12})$ of K α L α -shell (triplet state) and K β L α -shell (singlet state) for Li-atom

Fermi hole is plotted as a difference between $f(r_{12})$ of K α L α -shell (triplet state) and K β L α -shell (singlet state) for Li, by using data(2004) for mulit-term-Zeta, and both SZ,DZ for data (1974) by Clementi and Rotti presented in Figures3. From these figures it can be well-known in general that the radius of Fermi hole agreement in both SZ(0.013 when $r_{12} = 1.57$) and mulit-term Zate(0.011 when $r_{12} = 1.56$), compared this result with Ref [21] we found agreement between them, while Double-Zate given error in result presented in figure 3.

Reference

- [1] S,M.Mekelleche , A.B.Ahmed., J.Theor ChemAcc ,V103(2000).
- [2] I.I.Guseinov, B.A.Mamedov, J.Mathamatical Chemistry, V43(2007).
- [3] T.Koga, H.Matsuyama ,J. Chemical physics latters, V392(2004).
- [4] P. Maldonado, A. Sarsa, E. Bendî a, F.J. Gălvez, J. Atomic Data and Nuclear Data tables (2010).
- [5] C.C.Roothaan, J.Rev.Mod .phys, V32(1960).
- [6] N.Clementi and C.Roetti, J. AtomicData and NuclearData tables(1974).
- [7] A.Sarsa, F.J.Gălvez, E. Bendî a, J.Atomic Data and Nuclear Data tables(2004).
- [8] A.Sarsa, F.J.Gălvez, E. Bendî a, J. Phy.B.V36(2003).
- [9]M.L.Rustg and P.Tiwar, J. chemical; physics, V39(1963).
- [10]R.R.ZopeandB.I.Dunlap, j.chemical physics, V124(2006).
- [11]I.Guseionv,J.Mod.Model.v10(2000)
- [12]S.M.Mekelleche and A.B.Ahemd , J.Theor.chem.Acc,V103(2000).
- [14]K.E.Banyard and G.J.Seddon, V.58 (1973).
- [15]A.B.KarramPh.D Thesis college of education (Ibn-AL-Haitham) Baghdad (2004).
- [16] C.Bunge, J.Barrientas and A.Bunge, J. Atomic Data and Nuclear Data

Tables53(1993).

- [17] K.E.Banyard and K.H.AL-Bayati, J. Phys B; Atom Molec.phys., 19(1985).
- [18] M.N.Murshed, Ph.D. thesis college of education (Ibn-AL-Haitham) Baghdad (2001).
- [19] S.M.Ali, Ph.D, thesis college of science for women Baghdad university (2006).
- [20] N.H.Beebe and S.Lunell, J. Phys B; Atom Molec.phys.V8(1975),
- [21]E.M.AL-Robayai, Ph.D, thesis college of science for women Baghdad university (2006).