# Analysis of the electron energy distribution function and its transport coefficient in SF<sub>6</sub>-CO<sub>2</sub> applied gas mixtuers

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تحليل دالة توزيع طاقة الالكترون ومعلمات الحشد في خليط من غازي سداسي

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#### المستخلص

تم استخدام معادلة الأنتقال الطاقي لبولتزمان لحساب معلمات انتقال الألكترونات و دالة التوزيع الطاقي للألكترونات في كل من  $CO_2$ ,  $SF_6$  النقي وخليطيهما. وقد تم استخراج معلمات الحشد الألكتروني المتمثل بمعامل الانتشار للاكترونات و متوسط طاقة الاكترونات ضمن المدى ( $^{-10}$  الا $\times 5 \ge N / 3 \ge ^{-10}$  الانتشار للاكترونات و متوسط طاقة الاكترونات ضمن المدى ( $^{-10}$  الا $\times 5 \ge N / 3 \ge ^{-10}$  الانتشار الاكترونات و متوسط طاقة الاكترونات ضمن المدى ( $^{-10}$  الا $\times 5 \ge N / 3 \ge ^{-10}$  الانتشار الاكترونات و متوسط طاقة الاكترونات ضمن المدى ( $^{-10}$  الاكترونات و متوسط طاقة الاكترونات ضمن المدى ( $^{-10}$  الاكترونات و متوسط طاقة الاكترونات ضمن المدى ( $^{-10}$  الاكترونات و معامات الحشد الألكترونات و متوسط طاقة الاكترونات ضما المدى ( $^{-10}$  المحة الاكترونات و متوسط طاقة الاكترونات ضما المدى ( $^{-10}$  المحة الاكترونات و متوسط طاقة الاكترونات ضما المدى ( $^{-10}$  اللاكترونات و متوسط طاقة الاكترونات ضما المدى ( $^{-10}$  المحة الاكترونات و متوسط طاقة الاكترونات ضما المدى ( $^{-10}$  المحة الاكترونات و متوسط طاقة الاكترونات ضما المدى ( $^{-10}$  المحة مع معما المدى ( $^{-10}$  المحة مع ما المحة الما المحة الما المحة الما عالي المحة عنها معاد المحة الما عالي الما علي المحة والما معاد الالاتحة و المحة والمعام والمتمثلة في البرنامج المحة المحة الاكترون المحة و محاليات معاد المحة المحة المحة المحة مع الما معاد المحة المحة و محما المحة و المحة المحة و معاليات معاد المحة المحة المحة المحة و معاليات المحة المحة المحة الحقا المحة الحقال المحة و معاليات و محما لمحة المحة المحة المحة المحة المحة المحة الحقا المحة الحقا معاد المحة الحقا المحة و معاليات و المحة المحة و محة المحة الحقا المحة الحقا محة المحة المحة المحة المحة المحة المحة المحة الحقا محة المحة الحقا و محة المحة و محة المحة المحة المحة المحة المحة المحة و محة المحة الحقا و محة الحقا و محة المحة و محة الحقا و محة الحقا و محة المحة المحة الحقا و محة المحة الحقا و محة المحة الحقا و محة الحقا و

#### Abstract

The Boltzmann transport equation is used to calculate the electron energy distribution function (EEDF) and the transport coefficient in pure sulfur hexafluoride (SF<sub>6</sub>) and pure Carbon dioxide (CO<sub>2</sub>) and their mixtures. The electron swarm parameters are evaluated in the range of  $(5 \times 10^{-16} \le E / N \le 9 \times 10^{-15})$  V.cm<sup>2</sup>. These parameters namely are: The Diffusion coefficient of electrons and mean electron energy. The motion of electrons in plasma gas (SF<sub>6</sub>) and mixing it with (CO<sub>2</sub>) under an applied uniform electric field is simulated by using the numerical solution of Boltzmann's transport equation technique. The numerical solutions are utilized within the international computer code kinema-Elendif and written in Fortran 77 language software. The calculated distribution function is found to be remarked non-Maxwillian that has energy variations which reflect the import electron-molecule energy exchange processes.

**Keyword**: *Plasma, Electron Discharges, Dielectric property, Swarm Parameter, Kinetic Theory of Gases, Electron Energy Distribution Function.* 

# Introduction

In a matter of fact,  $CO_2$  cannot completely replace  $SF_6$  due to its low current interruption performance of high-voltage levels compared to that of pure  $SF_6$  gas. As a result, the most promising candidates to substitute  $SF_6$  gas is the  $SF_6$ - $CO_2$ mixtures. [3, 6]

A study of the electron energy distribution function in pure  $SF_6$  and  $CO_2$  and their mixtures are presented using the numerical solution of the Boltzmann equation for the Electron Energy Distribution Function in low-ionized plasma by using the code of kinema-Elendif which has been written in Fortran 77 language software. It is used to calculate the electron transport and kinetic coefficients in gas mixtures [7].

The data of momentum transfer (elastic collision) cross-section, the electronic excitation cross-section and the vibration cross-section for  $SF_6$  gas has been taken from Diefenbacher [8].

Three types of inelastic cross-section in  $CO_2$  have considered as:

- The vibration cross-section has been divided into four main vibration levels as given by Pitchford & Phelps [9] with onset energies (0.0827, 0.291, 0.580 and 0.870) eV.
- 2- The total ionization crosssection having onset energies of 13.3eV is given by Soonja *et al* [10].

For many years, the electric power industry has been using Sulfur Hexafluoride (SF<sub>6</sub>) gas as a dielectric and insulating material. In the event that replacement gases are consider a reasonable approach to reduce the use of SF<sub>6</sub> in high voltage electrical equipment.

The SF<sub>6</sub> gas is popular due to its unique physical and electrical properties as: nontoxic, nonflammable, noncorrosive, chemically stable with high breakdown strength and its dielectric strength is twice that of air [1, 2 and 3].

Considering the relatively poor dielectric strength of environmentfriendly pure gases and gas mixtures such as air, A gas used as a dielectric medium must have high dielectric strength, which is possible only with strong electronegative gases such as SF<sub>6</sub>, from a practical view, at partial concentrations of a few percent as possible substitute gases for pure SF<sub>6</sub>, which can reduce the negative effects atmosphere environmental on the pollution effectively. The candidates for pure gases mixed with  $SF_6$  are required to have sufficient insulation and current interruption capability [4, 5].

Eventually, the  $N_2$  and  $CO_2$  gases are the possible candidates that can be used for their chemical stability and no flammability or explosiveness as a result of gaining superior dielectric strength. In point of fact, the  $CO_2$  gas has started to gain attention as an arcquenching medium.

3- Excitation cross-section is given by Sierra et al [11].

#### Theory

#### **Boltzmann equation**.

important swarm parameters could be derived that it is still being used in many contemporary research projects to model transport phenomena [3].

The general form of the Boltzmann equation is [1, 2 and 3]:

The classical theory of transport processes is based on the Boltzmann transport equation; this equation can be driven simply by defining a distribution function and inspecting its time derivative. From this equation many

integral which accounts for electron energy transferred in elastic and inelastic collision.

The left hand side of the equation describe the behavior changes of electrons energy distribution function (EEDF) by the verity independent collisions while the right hand side describes the binary collisions of charges particles with the neutral gas species [5,6].

energy distribution function one can calculate the swarm parameters which they are diffusion coefficient of electrons and mean electron energy.

The diffusion coefficient of electrons is:-

Where f(r,v,t) is the electron velocity distribution function (EVDF) at time *t* and spatial location *r*, *V* is the electron velocity,  $\nabla_r$  is the gradient in *r*-space while (*e/m*) is the ratio of electron charge to its mass which is refers to the acceleration due to applied electric field (*E*) in (V.cm<sup>-1</sup>),  $\nabla_v$  is the gradient in V-

space and  $\left(\frac{\delta}{\delta t}\right)_{coll}$  is the collision

It is well known that the swarm parameters of electrons and collision cross-sections with molecules are related to each other's through the medium of the velocity distribution function of the swarm.From the electron

$$D = \frac{1}{3} \sqrt{\frac{2e}{m}} \int_{0}^{\infty} \frac{u}{N\sigma_{m}} f_{o}(u) d\varepsilon \qquad (2)$$

The Mean electron energy is [2]:

Where *N* is the gas density,  $\sigma_m$  is the momentum transfer cross section for elastic scattering and *u* is the electron energy expressed in electron volts.

$$\varepsilon = \frac{2}{3} \int_0^\infty u^{3/2} f_o(u, E / N, T) du \qquad (3)$$

Values of  $f_o(u)$  is calculated from Boltzmann's equation using overall collision cross-sections.

the electron transport or swarm parameters [4, 7].

#### **Calculations and results**

The EEDF for pure  $SF_6$  and pure  $CO_2$ and their mixture ( $SF_6$ - $CO_2$ ) for the different concentration are plotted as a function of E/N values and they are completely described in Figure (1, 2 & 3). It's necessary to note that there exists another mathematical technique for solving the Boltzmann transport equation by using the Monte Carlo method and involving the calculation of

The present study has resulted in a set of cross-sections for  $SF_6$  and  $CO_2$ which is consistent with measured swarm parameters for pure  $SF_6$  and  $CO_2$ respectively. The Reliability of these cross-sections and the Boltzmann equation procedure has been further tested by the comparison of measured and predicted values.



Figure (1): The distribution function as a function of E/N for pure  $SF_6$ .



Figure (2): The distribution function as function of E/N in pure  $CO_2$ .



Figure (3): The distribution function of  $(SF_6-CO_2)$  (50-50) % as a function of E/N value.

respectively for pure  $SF_6$ , pure  $CO_2$  and their mixtures as a function of E/N values.

The calculated transport coefficient which are the diffusion coefficient of the electron (D) and the mean electron energy ( $\varepsilon$ ) presented in figure (4 and 5)



Fig (4): The diffusion coefficient of as a function of E/N for different ratio gas mixtures.



Figure (5): The mean electron energy of as a function of E/N for different ratio gas mixtures

## **Discussion and conclusion**

collisions between electrons would cause the electron energy distribution function to tend towards a Maxwellian distribution. The influence of these collisions mainly depends on the ionization degree, but in Fig (3), the calculated distribution function are again markedly non-Maxwillian and having distinct varying curvatures at all electron energies. The pronounced dip distribution function in the are emphasized at low electron energy occur as a result of the high crosssections for vibrational excitations of (SF<sub>6</sub>, CO<sub>2</sub>) (50:50) mixture.

In Figurers (4, 5) which represents the diffusion coefficient and the mean energy of the electron respectively, in these two figure, the pure SF<sub>6</sub> has the values highest of the diffusion coefficient and mean electron energy while the pure  $CO_2$  has the lowest values of each of them. So, with decreasing SF<sub>6</sub> concentration in the mixture both of the diffusion coefficient and the mean electron energy is decrease and opposite of that with  $CO_2$ concentration decreasing the diffusion coefficient and the mean energy is increase this is implies that the number of collision in  $SF_6$  is lower than CO<sub>2</sub> as a result of the sets of crosssection data of the electron energy distribution in SF<sub>6</sub> and CO<sub>2</sub> gases and because of the growth of the inelastic collisions of electrons.

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It is important to say that the calculations of the electron transport parameters in  $(SF_6-CO_2)$  gas mixture reflect the advantages and disadvantages of each concentration mixing ratio. Therefore, one can choose the optimum cases that give a good compatible applied gaseous mixture.

A set of electron collision cross section, which is chosen to be as a consistent as possible with experiment to each of pure gases has been used. In Figure (1), the electron energy distribution function is affected by changing the parameter E/N. At low E/N values, the  $SF_6$  gas having higher attachment cross sections at low energy. This mean that, low E/N fluctuation results from the acceleration of each electron during the interval between collisions in the low energy region where the direction of the electrons is easily changeable by the electric field. Therefore, at high E/N values the  $SF_6$ become higher electronegativity gas. So, one can see the distribution functions are clearly non-Maxwillian and having variable distinct curvatures at all electron energies.

In figure (2), the vibrational energy losses in  $CO_2$  is distributing more energy, therefore, structure in f(u)is less apparent. Nevertheless, the dip in f(u) can be identified with electron energy loss to the asymmetric stretch vibration in  $CO_2$  for which the cross section is relatively large in this energy range. Also the distribution functions is clearly non- MaxwillianGenerally,

I wish to acknowledge the invaluable help of my dear husband

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