

Modalling Vibration of Dynamic Behavior of Power Plant Turbine

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Abstract

The paper explores the vibration of the turning shafts, and whirling hypothetically by managing just the pole bowing without torsion. It is demonstrated that the pole spinning results from vibration considered for this specific framework will have an essential whirling speed, which increments as the inactivity powers of the rotor. Whirling is related to quick pivoting shafts. When a pole pivots, it is subjected to outspread or radiating powers, which redirect the pole from its rest position. These radiating powers are unavoidable since the material inhomogeneities and get-together troubles guarantee that the focal point of gravity of the pole or its joined masses cannot harmonize with the pivot of revolution. The outward powers included and established that the main destabilizing or reestablishing power was because of the flexible properties or solidness of the pole. Subsequently, finding the essential speed, I made an endless redirection of the shaft due to whirling.

A stacking reaction investigation can reenact the reaction of the rotor by utilizing the Ansys program, which can distinguish fundamental modular recurrence. The modular investigation can give a serial of Eigen recurrence to a turning shaft, yet only some of them are the correct modular recurrence esteem we require. The consonant examination can separate the eigenvector of the modular investigation result, coordinate them together, and demonstrate the sufficiency recurrence reaction, twisting and stress dispersion qualities of some uncommon places of the turning shaft.

Keywords Turning shafts, Whirling, In homogeneities structure, Critical speed

Nomenclature

ω	Excitation	[rad/s]
E	Eccentricity	[mm]
K	Stiffness	[N/m]
Z	Damping ratio	[dimensionless]
C	Damping coefficient	[N.s/m]
M	Mass	[kg]
$X(\omega)$	Amplitude in X-direction	[mm]
$Y(\omega)$	Amplitude in Y-direction	[mm]
$G_x(i\omega)$	spectrum in X-direction	[N ² /Hz]
$G_y(i\omega)$	Spectrum in Y-direction	[N ² /Hz]
Φ	Phase angle	[degree]
a_{ij}	Deflections	[mm]
W_c	Critical speeds	[rad/s]

1.Introduction

In the present-day outline of turning hardware, the proportion of intensity created per pound of the pivoting components has quickly expanded. When planning the pivoting apparatus, the solidness conduct and the reverberation reaction can be obtained by figuring out complex Eigen esteems. The examination thinks about the impacts of immediate and cross-coupled coefficients of firmness and damping in the bearing and seals. Numerous papers have talked about how to enhance the edge execution of pivoting hardware, for example, changing the seal configuration, expanding the pole solidness, or receiving a more steady direction. All in all, any parameter change will influence the common recurrence and mode shape. Consequently, a reasonable decision of rotor firmness and mass conveyance may adequately enhance the strength of a rotor framework [1].

When a pole pivots, it goes into transverse motions. If the pole is out of adjustment, the subsequent radiating power will prompt the pole to vibrate when it turns at a speed equivalent to characteristic recurrence transverse motions. This vibration turns out to be huge and appears as a pole spinning. It likewise happens at products of full speed. This can harm substantial turning machines; for example, turbine generator sets and the framework must be precisely adjusted to decrease this impact, and they are intended to have a characteristic recurrence that is diverse to the speed of pivot. To begin or cease such apparatus, the essential speed must be kept away from it to prevent harm to the bearing and turbine cutting edges. Think about a weightless shaft [2].



Fig. (1) Rotating shaft of power plant

Vibration issues can happen whenever an engine is established or activated. When they happen, it is typically vital that one responds rapidly to take care of the issue. If settled slowly, we could either anticipate long-term harm to the engine or quick disappointment, which would bring about prompt generation loss. The loss of creation is the most basic worry; for this to happen, one should first comprehend the underlying driver of the vibration [3],[4],[5]. To take care of a vibration issue, one must separate amongst circumstances and results at the end of the day: Where does the power originate from the vibratory power, the reason for the large amounts of vibration, or is there a reverberation that enhances the vibratory reaction? Maybe the help structure is not sufficiently hardened to limit the dislodging [6],[7].

In this paper, the various sources of electrical and mechanical forces will be explained in the turbine Located in the Nasiriyah thermal station in the province of Thi Qar. Energy efficiency is essential to the economy as energy demand has continued to rise continuously over the past years. Since steam turbines are costly equipment, another solution must be found for replacement or re-maintenance. Components of high- and medium-pressure turbines are responsible for the highest loss of power output in heat plants and cause system failure in turbines. In order to maintain the state of the system, one must know the remaining life of the thermal turbines; this requires knowledge of the basic information of the system and the operating data (temperature, vibration, pressure, number and type of start, quality of fluid used, and data on work and irregularities), then find out the situation and as a result safety recommendations are put in place to continue working within a limited period, replacing or repairing some components in the event of a malfunction. The faults appear after 30 years of work, where efficiency starts less than 5%.

1.1 Dynamic Analysis of Rotor-Bearing System

The dynamic trademark for the rotor-bearing framework with a turning speed is entirely unexpected from the framework without pivoting. The fundamental motivation behind rotor dynamic examination is to decide the rotor-bearing framework's characteristic frequencies, mode shapes, basic speeds and unfaltering state reaction [8].

1.2 Rotor Unbalance

Static unbalance - All unequal masses lie in a solitary plane. This sort of imbalance can be recognized without turning the wheel .

1.3 Dynamic Unbalance

The unbalance is conveyed over different planes. This results in power and a shaking minute. The power can be identified as a static imbalance. However, the shaking minute must be controlled by turning [9].

2. Theoretical technique

The paper analyses a turbine located in Nasiriyah thermal station, Russian-made type K 210-130; it belongs to the older generation and is composed of high, medium and low pressure. The turbine's elements are chromium, vanadium, molybdenum and steel. This type produces 210MW and a pressure of 130 KG for any steam. It reaches the pressure of the steam turbine K 210-130 thermal efficiency, which is about 44.7%. It has been operating at Thi Qar Station since 1980. The analysis of the turbine in the ANSYS program is compared with the results taken from the work site to find solutions to problems that occur in the workplace. The most critical problems are journal bearing high clearance, parallel misalignment, bent shaft with a bend near the shaft centre, loose internal assembly and angular misalignment. The manufacturer's data are given in Table 1.

Table 1: Basic information about the HP turbine.

Basic information	Value
Fresh steam pressure inlet of stop valves, MPa	13
Rotor speed, rpm	3000
Design steam flow at 210 MW, t/h	628
Fresh steam temperature inlet of stop valves, °C	540
Outlet steam pressure of HP turbine, for design flow, Mpa	2.78
Maximum steam flow through a turbine, t/h	655
Relevant casing wall thickness, mm	90
Outlet steam temperature of HP turbine, °C	333



Fig. (2) Turbine axis

2.1 Whirling of rotating shafts

Fig. (3) Shows the displayed machines schematically, where the mass is amassed in a circle midway on the pole. Along these lines, the two degrees of opportunity are the translational removal of the circle. Both an adaptable shaft and spring bolsters are incorporated, and since the two segments have immaterial mass, their solidness might be joined utilizing the equation for springs in the arrangement. Finally, the backings will be viewed as highly adaptable to seclude vibration, in which case the rotor will be unbending [10]. The pole turns with the precise speed ω . Numerous mechanical frameworks include an overwhelming pivoting circle and a rotor appended to an adaptable shaft mounted on orientation. Run-of-the-mill cases are electric engines, turbines, blowers, and so forth. On the off chance that the rotor has some capriciousness, i.e., if the mass focus of the circle dosage does not correspond with the geometric focus, at that point, the turn delivers a diffusive power, making the pole twist. The revolution of the plane containing the bowed shaft about the bearing hub is known as spinning. For certain rotational speeds, the framework encounters brutal vibration and a marvel we propose to explore [11],[12].

Fig. (4a) demonstrates a pole pivoting with the consistent, precise speed ω concerning the inertial tomahawks x and y ; the pole conveys a circle of aggregate mass m at midcap and is considered massless. Consequently, the movement of the framework can be depicted by the removals x and y of the geometric focus S of the plate; even though this infers a two-level-of-opportunity framework, the x and y movements are autonomous, so the arrangement can be completed as though there were two frameworks with one level of flexibility each. As fundamental to the induction of movement conditions, we wish to determine how to speed up the mass focus. To this end, we indicate the inception of the inertial framework x and y by and the focal point of mass of the circle by C .

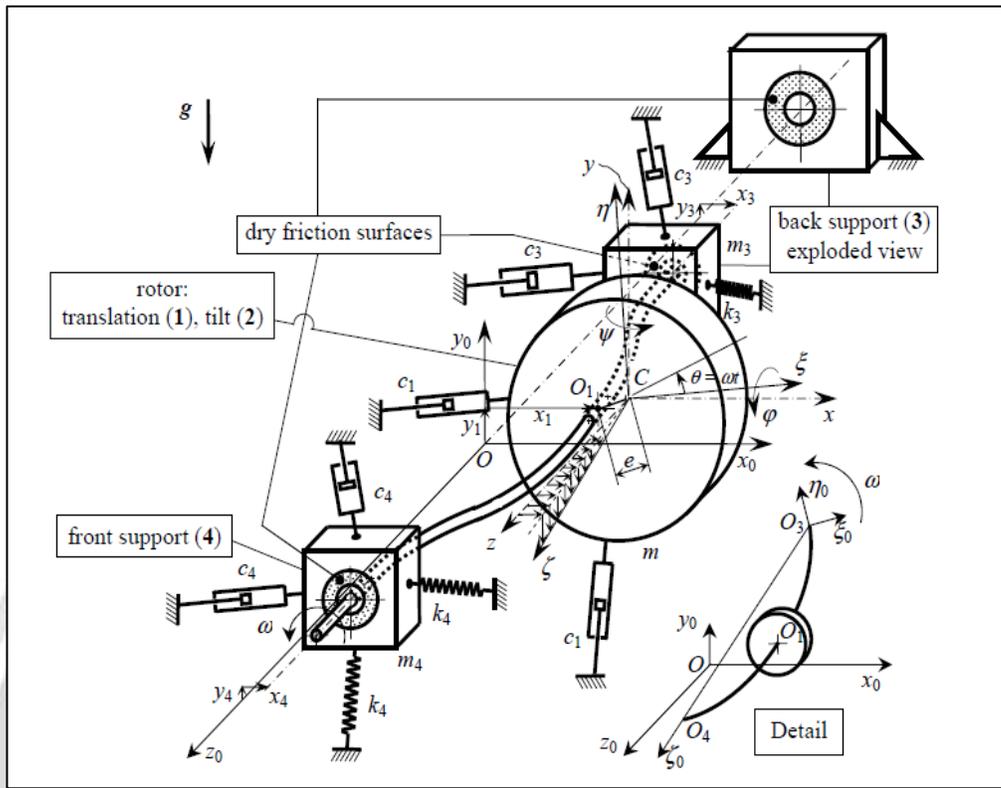


Fig. (3) Schematic of rotor bearing

Because of some rotor defect, the mass focus C does not agree with geometric focus S. We mean the separation between S and C by e, as in Fig (4b), where e speaks to the unusualness. To ascertain the quickening air conditioning of the mass focus C, we initially compose the span vector rc from O to C as far as rectangular segments as:

$$rc = (x + e \cos \cos \omega t)i + (y + e \sin \sin \omega t)j \tag{1}$$

Where i and j are constant unit vectors along axes x and y, respectively then differentiating Equation (1) twice with respect to time, we obtain the acceleration of C in the form

$$ac = (\ddot{x} - e\omega^2 \cos \cos \omega t)i + (\ddot{y} - e\omega^2 \sin \sin \omega t)j \tag{2}$$

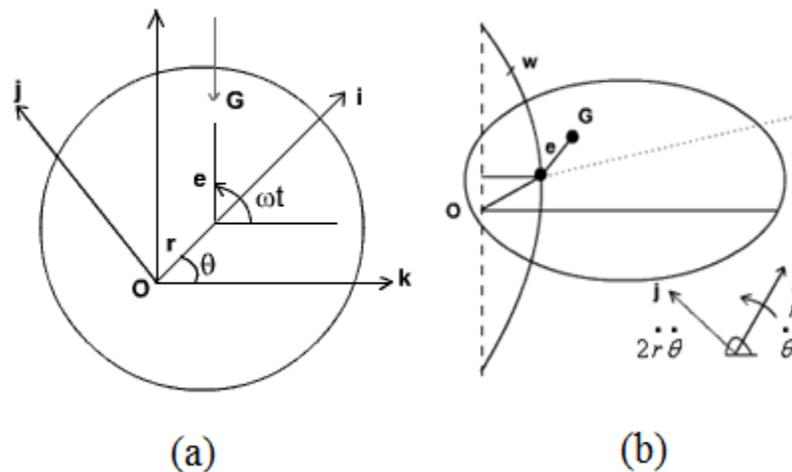


Fig. (4) Whirling of rotor bearing

To determine the conditions of movement, we accept that the main powers following up on the plate are reestablishing powers because of the versatility of the pole and opposing powers because of thick damping, for example, caused via air erosion [13].

The flexible impacts are spoken to by equal spring constants k_x and k_y , which are related to the distortion of the pole in the x and y bearings separately. Additionally, we accept that the coefficient of gooey damping is the same in the two bearings and equivalent to c . the flexibly reestablishing powers and the thick damping powers are acting at point s . thinking about Equation (1), the x and y part of Newton's second law, Equation (2) are

$$-k_x x - c\dot{x} = m(\ddot{x} - e\omega^2 \cos \cos \omega t) \quad (3)$$

$$-k_y y - c\dot{y} = m(\ddot{y} - e\omega^2 \sin \sin \omega t)$$

This can be rearranged in the form

$$\begin{aligned} \ddot{x} + 2\zeta_x \omega_n \dot{x} + \omega_{nx}^2 x &= e\omega^2 \cos \cos \omega t \\ \ddot{y} + 2\zeta_y \omega_n \dot{y} + \omega_{ny}^2 y &= e\omega^2 \sin \sin \omega t \end{aligned} \quad (4)$$

Where

$$\zeta_x = \frac{c}{2m\omega_{nx}}$$

$$\omega_{nx} = \sqrt{\frac{k_x}{m}} \quad (5)$$

$$\zeta_y = \frac{c}{2m\omega_{ny}}$$

$$\omega_{ny} = \sqrt{\frac{k_y}{m}} \quad (6)$$

are viscous damping factors and natural frequencies.

This should not shock anyone, as a pivoting uneven mass. Consequently, the unflinching state arrangement of condition (4) can be acquired, after that example, we can compose the arrangements

$$\begin{aligned} x(t) &= |X(\omega)| \cos(\omega t - \phi_x) \\ y(t) &= |Y(\omega)| \sin(\omega t - \phi_y) \end{aligned} \quad (7)$$

Where the individual amplitudes are

$$\begin{aligned} |X(\omega)| &= e^{\left(\frac{\omega}{\omega_{nx}}\right)^2} |G_x(i\omega)| \\ |Y(\omega)| &= e^{\left(\frac{\omega}{\omega_{ny}}\right)^2} |G_y(i\omega)| \end{aligned} \quad (8)$$

In which

$$\begin{aligned} |G_x(i\omega)| &= \frac{1}{\left\{ \left[1 - \left(\frac{\omega}{\omega_{nx}} \right)^2 \right]^2 + \left(\frac{2\zeta_x \omega}{\omega_{nx}} \right)^2 \right\}^{\frac{1}{2}}} \\ |G_y(i\omega)| &= \frac{1}{\left\{ \left[1 - \left(\frac{\omega}{\omega_{ny}} \right)^2 \right]^2 + \left(\frac{2\zeta_y \omega}{\omega_{ny}} \right)^2 \right\}^{\frac{1}{2}}} \end{aligned} \quad (9)$$

are magnitudes and

$$\begin{aligned} \phi_x &= \frac{\frac{2\zeta_x \omega}{\omega_{nx}}}{1 - \left(\frac{\omega}{\omega_{nx}} \right)^2} \\ \phi_y &= \frac{\frac{2\zeta_y \omega}{\omega_{ny}}}{1 - \left(\frac{\omega}{\omega_{ny}} \right)^2} \end{aligned} \quad (10)$$

are the phase angles

One consider first the most well-known case, in particular, that of a pole of roundabout cross area, so the solidness is the same in the two headings, $k_x = k_y = k$. For this situation, the two regular frequencies match thus do the thick damping factors, or

$$\omega_{nx} = \omega_{ny} = \omega = \sqrt{\frac{k}{m}} \quad (11)$$

Moreover, in view of Equation (11), we conclude from Equation (9) and (10) that the magnitudes on the one hand and the phase angles on the other hand are the same ,or

$$|G_x(i\omega)| = |G_y(i\omega)| = |G(i\omega)| = \frac{1}{\left\{ \left[1 - \left(\frac{\omega}{\omega_n} \right)^2 \right]^2 + \left(\frac{2\zeta\omega_n}{\omega} \right)^2 \right\}^{\frac{1}{2}}}$$

$$\phi_x = \phi_y = \phi = \frac{\frac{2\zeta\omega}{\omega_n}}{1 - \left(\frac{\omega}{\omega_n} \right)^2} \quad (13)$$

It follows immediately, from Equation (7), that the amplitudes of the motions x and y are equal to one another, or

$$|X(\omega)| = |Y(\omega)| = e \left(\frac{\omega}{\omega_n} \right)^2 |G(i\omega)| \quad (14)$$

But, from Fig (2) and Equation (7) we can write

$$\omega t - \phi) \quad (15)$$

From which we conclude that

$$\theta = \omega t - \phi \quad (16)$$

$$\text{and } \dot{\theta} = \omega \quad (17)$$

Henceforth, in this situation, the pole spins with an indistinguishable precise speed from the revolution of the circle, so the pole and the plate pivot together as an unbending body; this is known as no concurrent spin. It is anything but difficult to confirm that in synchronous spin, the spiral separation from O to S for a given ω is steady,

$$r_{os} = \sqrt{x^2 + y^2} = e \left(\frac{\omega}{\omega_n} \right)^2 |G(i\omega)| = \text{constant} \quad (18)$$

So, point S describes a circle about point O. To determine the position of C relative to the whirling plane; we consider Equation (16). The relation between the angles is depicted in Fig. (4). Indeed, from Fig. (4), we can interpret the phase angle as the angle between the radius vectors. Hence, recalling the second of Equation (13), one concludes that.

As a final remark concerning synchronous whirl, we note from Equation (13) that the magnitude and the phase angle have the same expressions as in the case of the rotating

unbalanced mass, which corroborates our earlier statements that the two systems are analogous [14]. Next, we return to the case; in this case, solution (8) can be written as

$$x(t) = X(\omega)\cos\omega t \quad (19)$$

$$\text{where } X(\omega) = \frac{e\left(\frac{\omega}{\omega_{nx}}\right)^2}{1 - \left(\frac{\omega}{\omega_{nx}}\right)^2}$$

$$Y(\omega) = \frac{e\left(\frac{\omega}{\omega_{ny}}\right)}{1 - \left(\frac{\omega}{\omega_{ny}}\right)} \quad (20)$$

Dividing the first of Equation (19) by $X(\omega)$ and the second by $Y(\omega)$, squaring and adding the results, we obtained

$$\frac{x^2}{X^2} + \frac{y^2}{Y^2} = 1 \quad (21)$$

Which represents the equation of an ellipse. Hence, as the shaft whirls, point S describes an ellipse with O as its geometric center. To gain more insight into the motion, we consider Equation (19) and write

$$\tan \theta = \frac{y}{x} = \frac{Y}{X} \tan \omega t \quad (22)$$

Differentiating both sides of Equation (22) with respect to time and considering Equation (19), we obtain

$$\dot{\theta} = \frac{XY}{X^2 \cos^2 \omega t + Y^2 \sin^2 \omega t} \omega \quad (23)$$

But the denominator on the right side of Equation (23) is always positive, so that the sign of $\dot{\theta}$ depends on the sign of XY . By convention, the sign of ω is assumed as positive, i.e., the disk rotates in the counter-clockwise sense. one can distinguish the following cases:

1. $\omega < \omega_{nx}$ and $\omega < \omega_{ny}$. In this case, one concludes from Equation (19) that $XY > 0$, so that point S moves on the ellipse in the same sense as the rotation ω .

2. $\omega_{nx} < \omega < \omega_{ny}$ or $\omega_{ny} < \omega < \omega_{nx}$. In either of these two cases $XY < 0$, so that S moves in the opposite sense.

3. $\omega > \omega_{nx}$ and $\omega > \omega_{ny}$. In this case $XY > 0$, so that S moves in the same sense as ω . Examining solution (18) and (19) for the undamped case, we conclude that the possibility of resonance exists. In fact, there are two frequencies for which resonance is possible, namely,

$\omega = \omega_{nx}$ and $\omega = \omega_{ny}$. Of course, in the case of resonance, solutions (18) and (19) are no longer valid [14]. It easy to verify by

Substitution that the particular solutions in the two cases of resonance are

$$\begin{aligned} x(t) &= \frac{1}{2} e \omega_{nx} t \sin \sin \omega_{nx} t \\ y(t) &= \frac{1}{2} e \omega_{ny} t t \end{aligned} \quad (24)$$

2.2 Whirling of shaft

Whirling is characterized as the turn of the plane made by the twisted shaft and the line of the focal point of the bearing. It happens because of various variables, some of which may incorporate (i) unconventionality, (ii) lopsided mass, (iii) gyroscopic powers, and (iv) liquid grinding in bearing thick damping..

2.3 Critical speeds analysis

The shaft will bend to the most straightforward shape possible at the first critical speed. At the second critical speed, it will be bent to the second simplest shape possible. .

2.4 Dunkerly equation

Where w_c is the first critical speed of multi mass system.

W_1 : is the critical speed which would exist if only mass no 1 where present.

Higher critical speed:

For multi mass system requires more exterior calculation than is necessary forth determine of the lowest .first critical speed [15].

$$\frac{1}{\omega^4} - (a_{11}m_1 + a_{22}m_2) \frac{1}{\omega^2} + (a_{11} a_{22} - a_{12}a_{21}) m_1 m_2 = 0 \quad (25)$$

$$\left| \begin{array}{cccc} a_{11}m_1 - \frac{1}{\omega^2} & a_{12}m_2 & a_{13}m_3 & a_{21}m_1 \\ a_{12}m_2 & a_{22}m_2 - \frac{1}{\omega^2} & a_{23}m_3 & a_{21}m_1 \\ a_{13}m_3 & a_{23}m_3 & a_{33}m_3 - \frac{1}{\omega^2} & 0 \\ a_{21}m_1 & a_{21}m_1 & 0 & a_{11}m_1 - \frac{1}{\omega^2} \end{array} \right| = 0 \quad (26)$$

3. Result and Discussion

Vibrations in the turbine are due to various issues, such as shortcomings of the rotating parts, unreasonable clatter, or vibration transmission to the supporting structure. An imperative wellspring of this vibration is out-of-balanced qualities, and this paper prescribes that the rotor response is decreased by suspending the machine on nonlinear springs and adding mass to lessen vibration and clamour. In vibration separation, nonlinear mounts have been proposed which have an indistinguishable static strength from an indistinguishable straight sponsorship, i.e. load, bearing capacity, and anyway, meanwhile, offer a low component robustness, i.e. a lower typical repeat. Like this, the isolator is excellent for extended repeat runs. These mounts are alluded to in

the composition as high-static-low-dynamic-immovability (HSLDS) frameworks. In this paper, the rotor is suspended on a setting HSLDS spring to broadly diminish the separating rates to values far from the working speed. The upsides of the nonlinear sponsorships are demonstrated using a precise two-level adaptability rotating machine show, including a rigid circle, shafts, course, and support that are versatile and have immaterial mass. After an immediate examination to feature the benefits of low component immovability, an expected symptomatic game plan of the nonlinear numerical articulation of development is displayed. A relationship between the straight and nonlinear response shows the reasonability of the nonlinear sponsorships. Finally, the issues that happen if the nonlinearity is too strong are featured.

3.1 Geometry

The turbine was drawn using the Sold Work program according to the dimensions of the Nasiriyah thermal station. As shown in Fig. (5), the mesh of the turbine structure is meshed by the program for nodes and elements, as shown in Fig. (6).

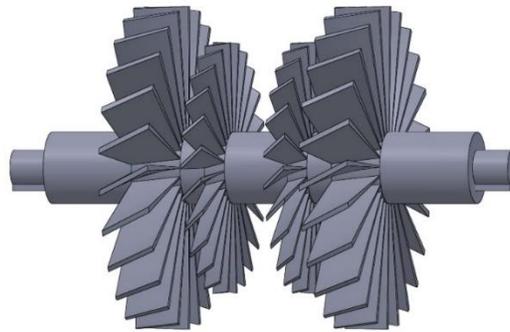


Fig. (5) Drawing in the Sold Work program

Statistics	
Nodes	29100
Elements	11526

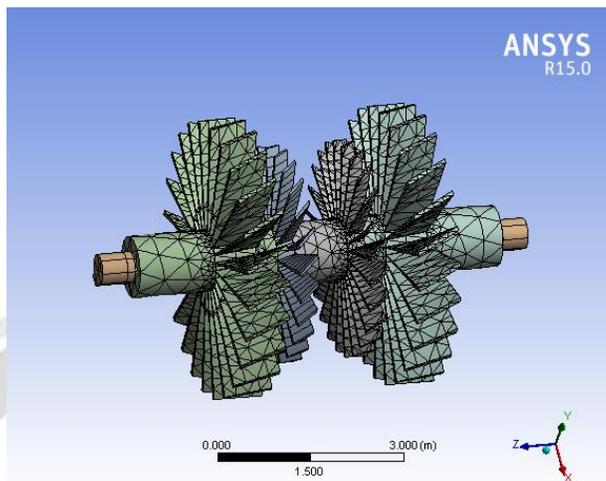


Fig. (6) Mesh of the structure

After the smearing moment, we appear for deformation and equivalent stress in the form as they are the warning factors for design. Figures. (7-8) show this structure's maximum deformation and maximum equivalent stress through static loading. The maximum deformation near the ends of the blades has a value of 6.4771 m , and maximum equivalent stress is also observed at the centre of the turbine and has a value of 2.0754 pa .

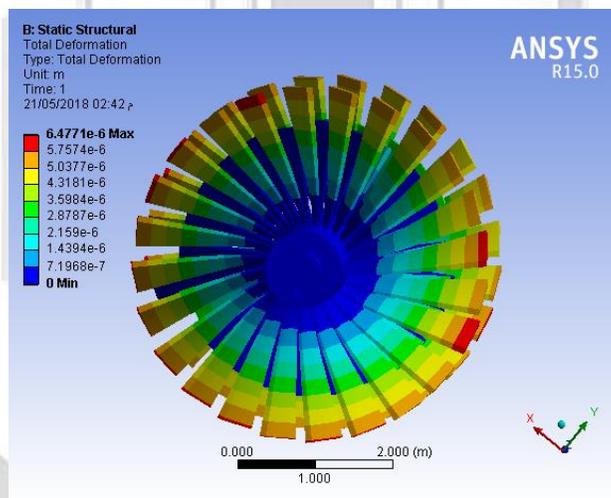


Fig. (7) Total deformation for static structural

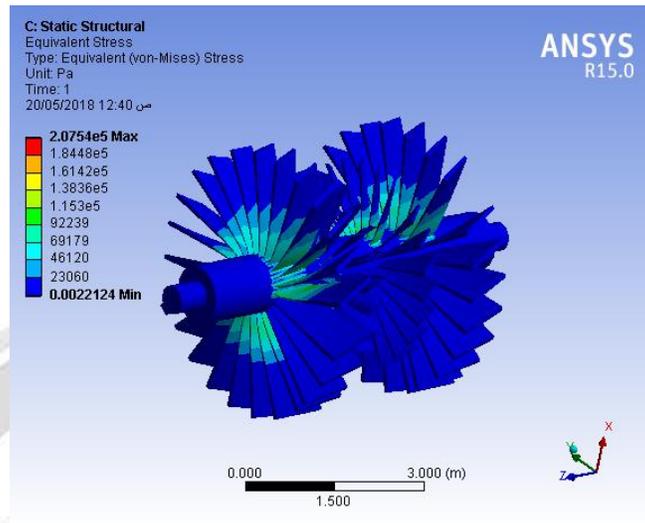


Fig. (8) Equivalent stress for static structural

3.2 Rotating Dynamic

Since turning mechanical assembly needs to rotate to do supportive work, we ought to consider the result for the primary strategy for our rotor once it is turning. However, again, we will have three particular versions with growing bearing immovability, and we will expect our support introduction to have approach strength in each winding heading. We should review our examination/particular test with the shaft turning at 10 rpm. We should look at the repeat and mode condition of the most insignificant ordinary repeat and the frequencies and mode shapes for a minimal technique for the three machines. Note that the condition of the development has changed. In any case, the frequencies are close to the no-turning first mode. As in the no-turning case, the bearing immovability to shaft strength extent significantly impacts the mode shape. However, again, the case with no shaft bowing is suggested as an unbendable mode. These modes look all that much like the no insurgency modes, yet they presently incorporate circuitous development rather than planar development.

First, imagine swinging a jump rope around to picture how the rotor is moving. The rope takes after the graph of a swelling barrel. In like manner, this mode is now and again suggested as a "round and empty" mode. Seen from the front, the restriction gives an impression of being swaying all over. Therefore, this mode is called a "sway" or "translator" mode a portion of the time. Not at all like most jump ropes, but then again, the rotor is turning. The turning development of the rotor (the 'bounce rope development') can be done in an indistinguishable bearing from the shaft's insurgency or another way. Rotor navigates the course of time for both synchronous forward and synchronous in switch turn. Note that for a forward turn, a point on the surface of the rotor moves in an indistinguishable bearing from the turn. In this way, for synchronous forward turn (unbalanced excitation, for example), a point at the outside of the rotor remains at the outside of the turn circle. In an inverted turn, on the other hand, a point at the surface of the rotor moves the other path as the turn to inside the turn hovers in the midst of the turn. To see how a broader extent of shaft speeds changes the situation, we could play out the examination/particular test with an extent of shaft paces from non-swinging to fast. We could

then take after the forward and turn-around frequencies associated with the primary mode. Figures 9 and 10 plot the Amplitude for Rotational and total translational amid associated minute unmistakable exact standard frequencies over a broad shaft rate run. The frequencies of this round and empty mode do not change significantly when finishing the speed run. The retrogressive turn mode drops imperceptibly, and the forward turn mode increases somewhat (most recognizably in the high-solidness case). The reason behind this change will be examined in the accompanying fragment.

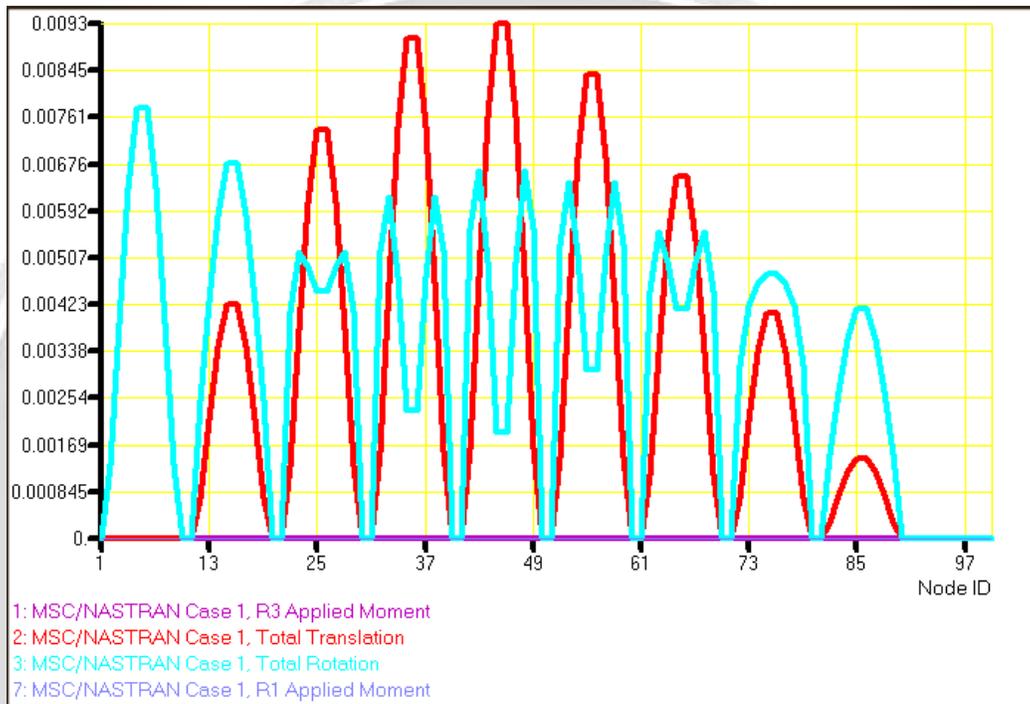


Fig. (9) Rotational and total translational during applied moment

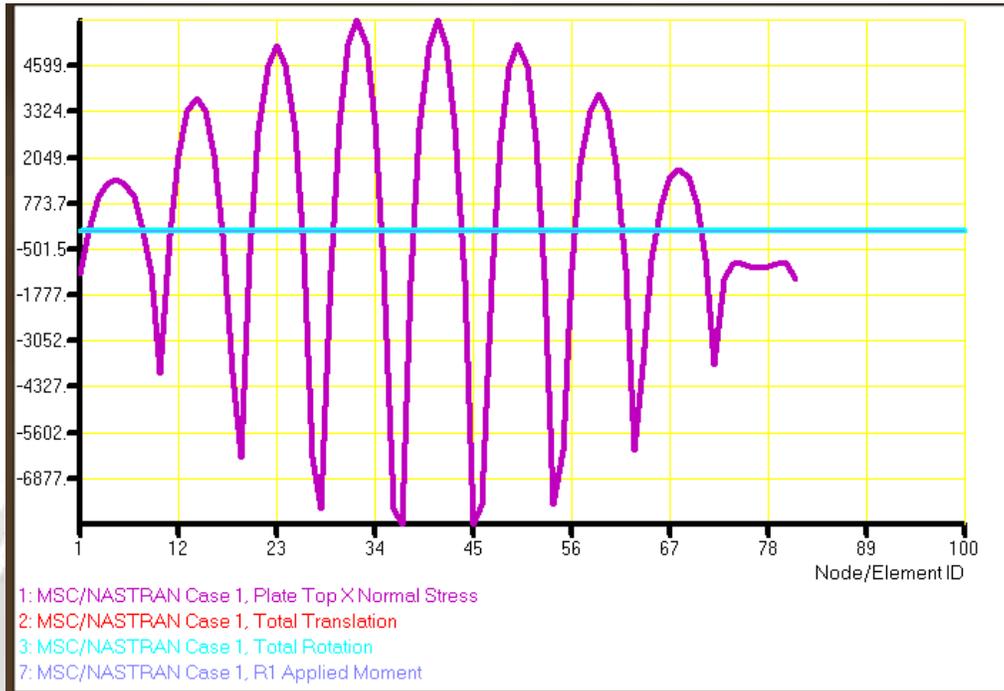


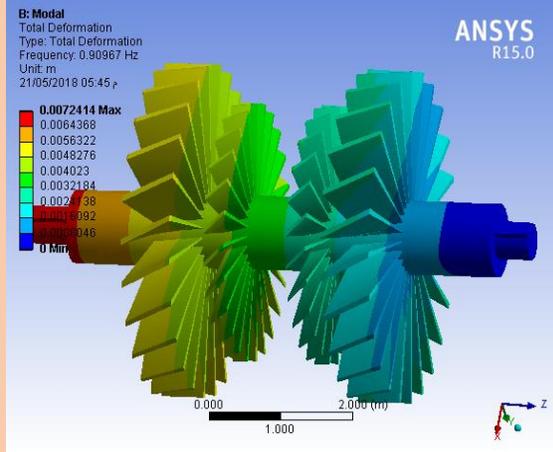
Fig. (10) Normal stresses in the rotating shafts.

3.2 Modal analysis

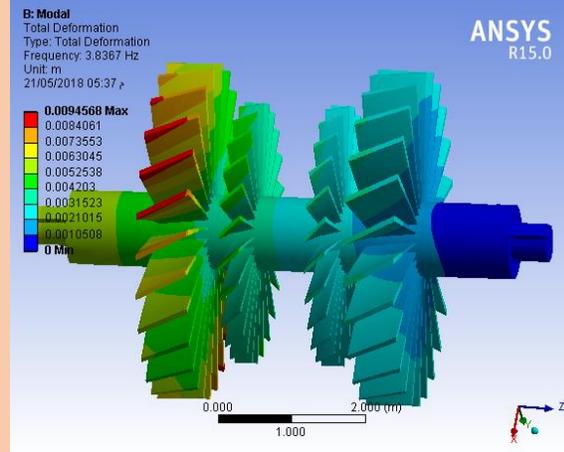
Most importantly, a modular investigation of this rotor is prepared. From the consequence of modular examination, modular recurrence and shape can be understood, which are helpful for the sensor dispersion configuration of attractive bearing. The first five mode shapes are shown in Fig. (11), and the first five natural frequencies are shown in table 2:

Table 2. Natural frequency of particular mode shapes

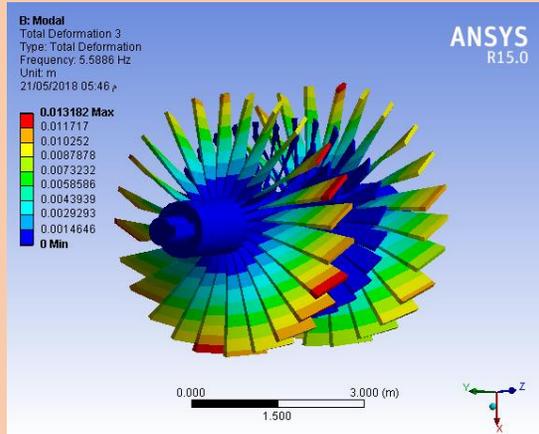
Mode	1.	2.	3.	4.	5.
Frequency[Hz]	0.90967	3.8367	5.5886	5.8003	5.8095



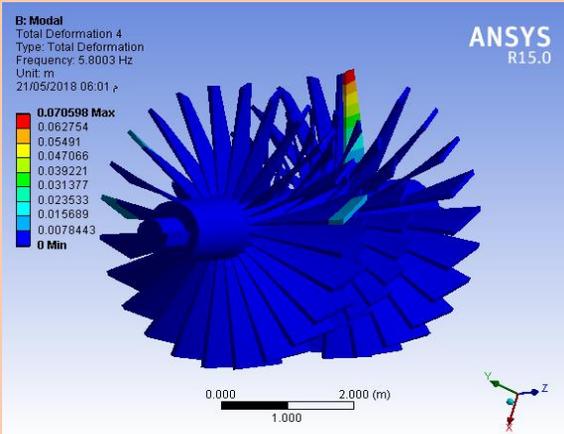
First mode shape at natural
frequency 0.90967Hz



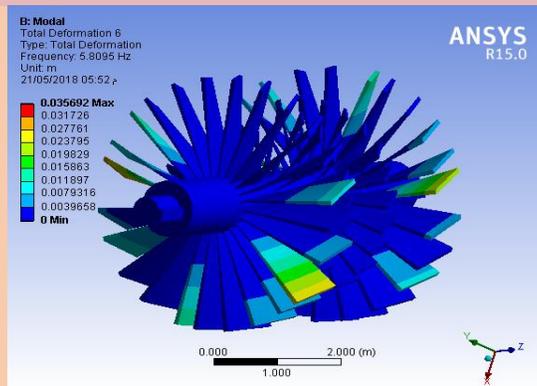
Second mode shape at natural
frequency 3.8367Hz



Third mode shape at natural
frequency 5.5886Hz



Fourth mode shape at natural
frequency 5.8003Hz



Fifth mode shape at natural frequency 5.8095Hz

Fig. (11) The first five mode shapes of power plant turbine.

The following result shown in Figures (12-14) is of the workplace vibration (Harmonic response). It is observable that there is very limited quantity of stress and deformation being induced in the body. The maximum equivalent stress is 15.543 Mpa whereas maximum deformation is 9.8516 mm and maximum equivalent strain 9.0029mm.

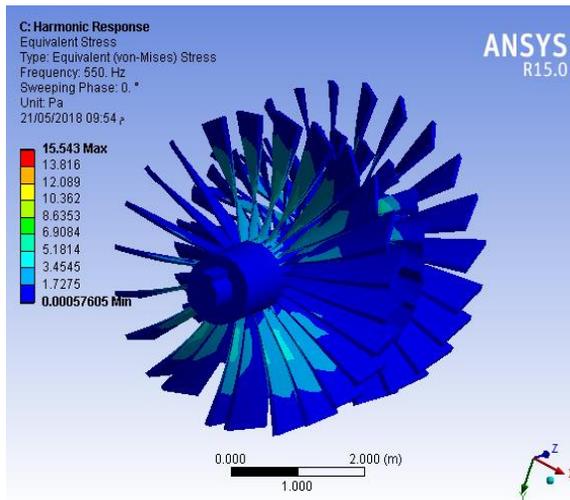


Fig. (12) Equivalent stresses harmonic response.

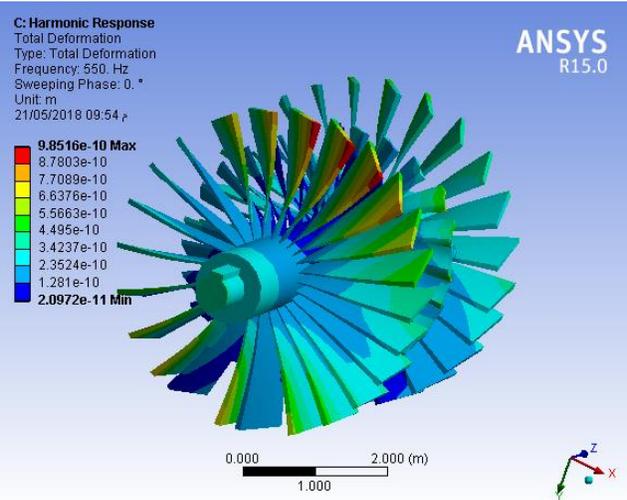


Fig. (13) Deformation for harmonic response.

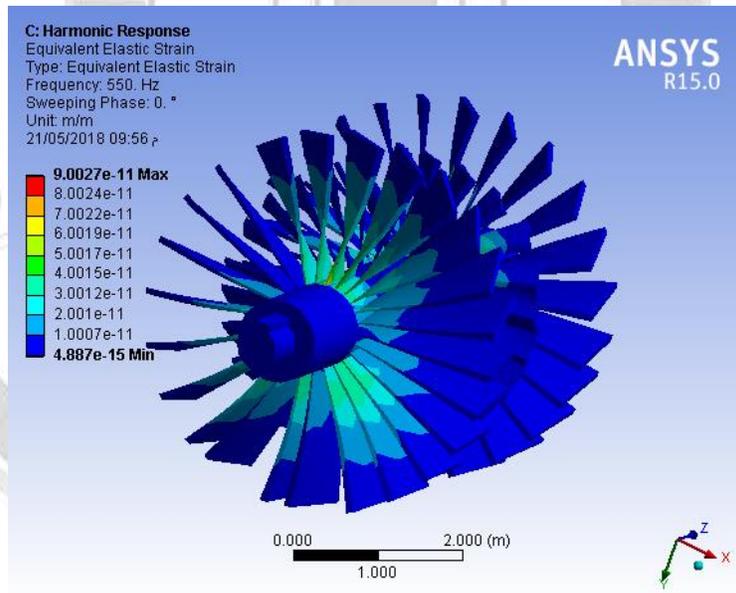


Fig. (14) Equivalent strian harmonic response.

Fig. (15) This represents the course of the turbine over time and illustrates the problems that occur at a given time. We note from the form that the system is balanced at the start of the station's operation, but note after some time that the vibration value is increased due to the reasons mentioned earlier so that these problems are treated with precise identification of the cause of the problem.

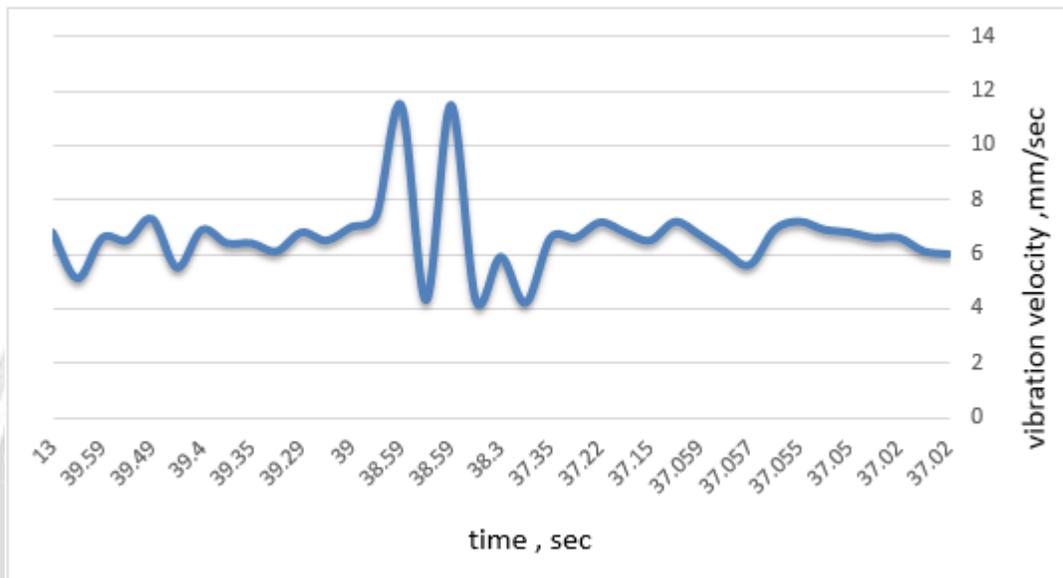


Fig. (15) Relation between vibration velocity and time.

To become Frequency Response Function (FRF) in ANSYS-15 Workbench. A mode superposition harmonic analysis is conducted. Figure (16 – 17) show the FRF for turbine without and with add a mass of balance and reduce vibration respectively. These figures illustration that the harmonic amplitudes were reduced in the modes with using add a mass of balance and reduce vibration.

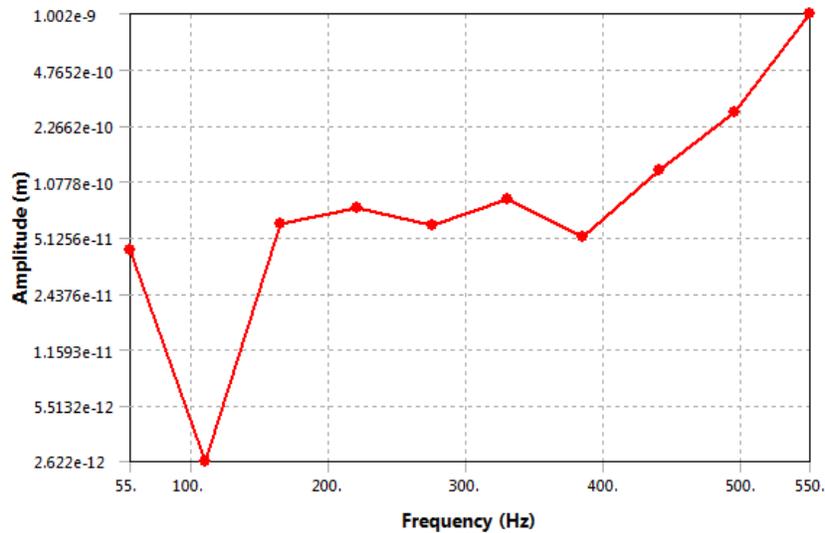


Fig. (16) FRF for turbine without add a mass of balance and reduce vibration.

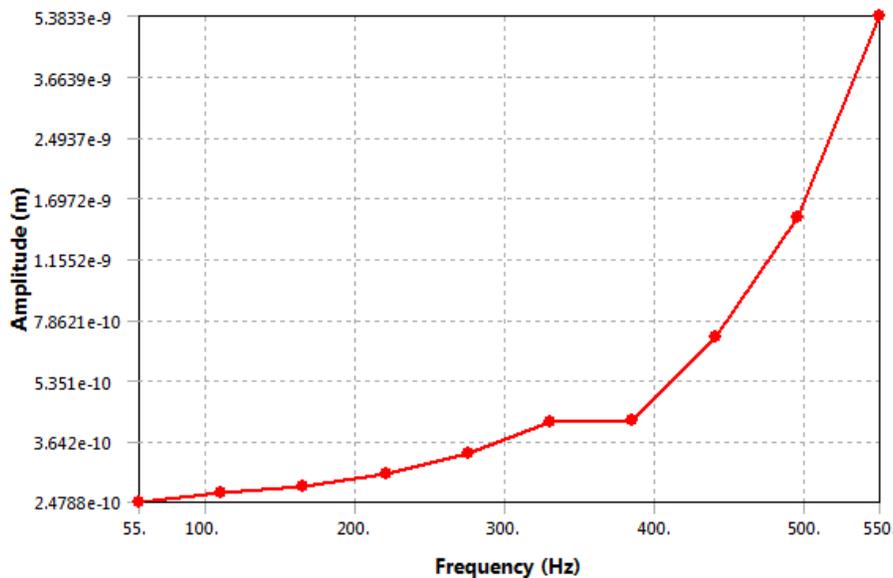


Fig. (17) FRF for turbine with add a mass of balance and reduce vibration.

4. Conclusion

In this paper, the dynamic behaviour of Turbine adaptability considers the hypothetical results and the numerical reenactment. The results are pretty acceptable compared to those obtained from Nasiriya Thermal Station. Our work is a solution to one of the problems located in Nasiriya Thermal Station (journal bearing high clearance, parallel misalignment, bent shaft with a bend near the shaft centre, loose internal assembly and angular misalignment). The work was

limited to the analysis of turbine blades and the defects that occur in them and methods of treatment to reduce the vibration and noise resulting from this imbalance and reduce the causes of problems as much as possible with the lowest cost. The methods used to deal with noise and vibration resulting from the lack of balance or whirling are the addition of a mass of the balance, the addition of insulation and others.

It was observed that the turning development under mechanical resonance pursues the turning speed, and it can be seen subjectively in the turning development. Exactly when the versatile rotor structure, like a flywheel, is related to a generator, the pursuing possibly appears in the electromotive intensity of the generator when the flywheel is under mechanical resonance. The change is made by adjusting the damping capacity of the rotor. In any case, the system is disadvantageous for the flywheel essentialness amassing structure because of the imperativeness hardship. The utilization of the structure for energy amassing is an extraordinary hazard when the versatility of the post appears. In the following stage, we will add to dynamic control of the turning development in the flywheel structure. The got gives us snippets of data to change the turning speed in the versatile rotor structure. Bearing cabin components on the response traits of a versatile rotor system with a partner flexibility bearing.

The dynamic arrangement of the rotor system is to upgrade the structure in quality at the working pace, unbalance response in the locale of the rotor essential speeds, and limit the system weight considering monetary issues. This paper deals with the perfect shape setup of the rotor shaft to change the separating speeds under the goals of the enduring mass. Regarding the setup strategies, the genetic count was associated with finding the perfect separations crosswise over a rotor shaft so the upgraded rotor structure can yield the essential speeds as far from the working velocity as expected in light of the current situation. The innate count is a chase computation considering the typical world, which communicates that the individual with the transcendent character influences the pervasive relative due to having splendid adaptability, high survivability, and all the more crossing point shot, and the normal inherited that on and on play out the strategy of age, crossover, and change. The results exhibit that the fundamental paces of the rotor-bearing structures can be upgraded by a slight change of the post-removal, even without growing the total mass of a rotor system.

Interest Statement

There is no reason to explain why the straight axes are not balanced with their vertical axes when they rotate at different speeds. Application and practice are found at certain speeds. The axis takes the deviation further, and the spin occurs for the axis.

If the speed of the axle is maintained at a critical speed, the deviation at the axis increases sufficiently to cause the axis to fail. Critical acceleration usually occurs when the spin appears and appears at any other speed. The axis allows access to equilibrium. If the axis speed increases sufficiently and quickly, it is possible to pass through the critical speed before an abnormal rotation displacement. To choose the balance of the mechanical system, it is possible to affect the state of balance and observe which follows the movement that tends to fade and leave the system more and faster at equilibrium or whichever tends to increase and take the system away from the state of equilibrium.

By adding more than one mass of the axis and increasing the critical velocity on the condition that the gyroscopic effects are minor. The critical velocity corresponds to the natural frequencies, and the axis will be vertically aligned to avoid any disturbance of the angular angle by considering the effects of gravity. The weight has no effect if the axis has a circular section, the axis deviates under the influence of its weight, and the spin takes place around the deviation of the midline. The occurrence of the critical speed of the entire turbine will lead to a break in the blades of the turbine. Vibration is unacceptable because it is a source of noise, and this vibration can generate a range of audible frequencies that produce a sound capable of harassing staff at the station.

I suggest improving the performance of the turbine by coating the turbine with materials to protect it from corrosion.

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الاهتزازات النموذجية للسلوك الديناميكي في توربينات محطات الطاقة

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الخلاصة

في هذا البحث تم دراسة اهتزازات وتشويش الاعمدة الدوارة المنحنية نظريا دون التطرق الى الاهتزازات اللاتوائية حيث ان التشويش في الاعمدة الناشى من الاهتزازات لهذا النوع الخاص من الانظمة يمتلك سرع تشويش حرجة تزداد بزيادة قوى القصور للاعمدة الدوارة. ان التشويش يقترن بالاعمدة الدوارة السريعة حيث عندما تدور الاعمدة فانها تتعرض الى قوى طاردة مركزية تسبب انحراف العمود بعيدا عن موقع الاستقرار الخاص به حيث ان هذه القوى الطاردة المركزية لايمكن تجنبها لان عدم التجانس في المادة وصعوبات التجميع تحيل دون تطابق مركز الثقل للعمود او مراكز ثقل الكتل المرتبطة به مع محور الدوران. ان القوة الطاردة المركزية تضمن وتحدد بان قوى الاسترجاع والتوازن هي تلك الناتجة عن الخواص المرنة او الصلابة للعمود ولهذا السبب نستنتج بان السرعة الحرجة هي التي تسبب انحراف غير محدد نتيجة التشويش. ان تحليل الاستجابة القسرية يمكن الاستفادة منها في محاكاة الاستجابة باستخدام برنامج Ansys ويمكن المساعدة لتحديد الترددات النسقية الرئيسية. ان التحليل النسقي بالنسبة للاعمدة الدوارة يعطي سلسلة من النسوق نحتاج الى البعض منها للدراسة في هذا المشروع. وان التحليل المتجانس يستخلص النسوق المطلوبة وباستخدام النسوق والترددات سوية يمكن معرفة خواص الاستجابة الترددية وتوزيع الاجهادات والتشوهات لمواقع معينة للعمود الدوار.

كلمات الدالة: الاعمدة الدوارة، التشويش، الهياكل الغير متجانسة، السرعة الحرجة.